

ATELIER GAINES PLASMAS
MARSEILLE, NOVEMBER 5TH 2024

END-ELECTRODE BIASING FOR ROTATION CONTROL: MODEL AND EXPERIMENTS IN LAPD

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▶ Motivations

- Why controlling plasma rotation ?

▶ End-electrode control: when should it work?

- Sheath
- Voltage drop along field lines
- Necessary conditions

▶ Experimental campaign on LAPD

- Setup
- Results

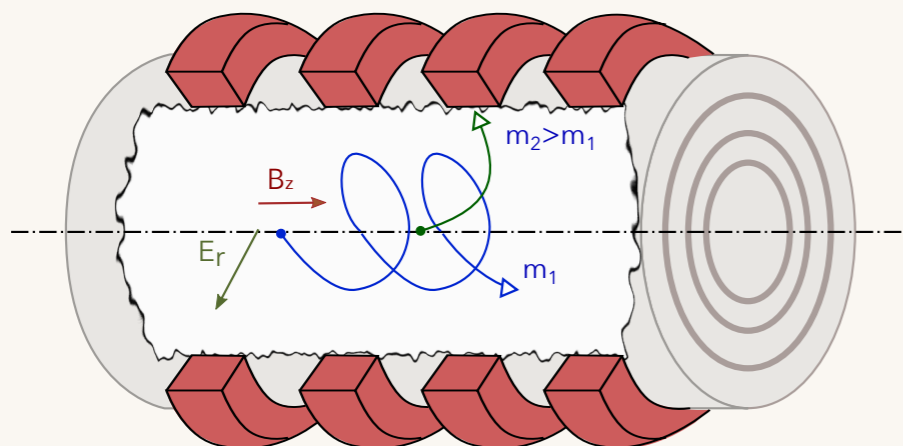
▶ Conclusion

WHY ROTATION CONTROL ?

Largely explored for instabilities and turbulence curbing or suppression. Also interesting in the confinement properties rotation creates.

Plasma mass separation

Leverage the differential confinement properties offered by a rotating plasma
[think e. g. of a plasma centrifuge]



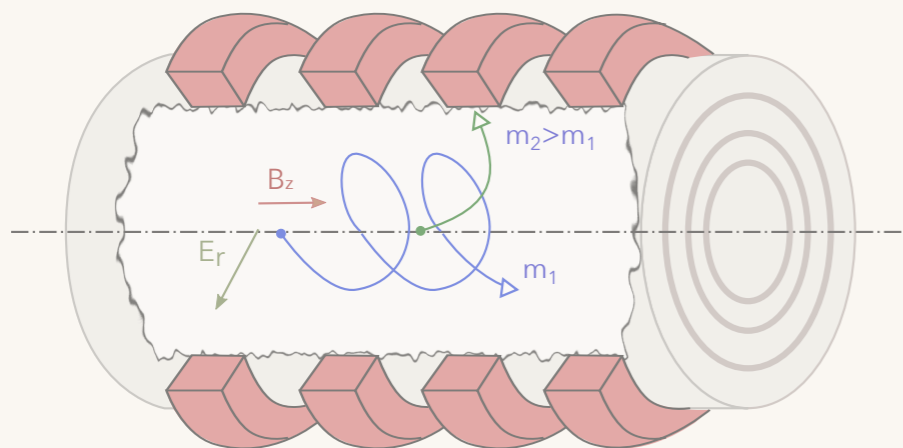
Technology attractive for nuclear fuel cycle,
nuclear clean up, rare earth recycling...

Zweben, Gueroult and Fisch (2018), Phys. Plasmas, **25**, 090901
Gueroult et al. (2018), PPCF, **60**, 014018

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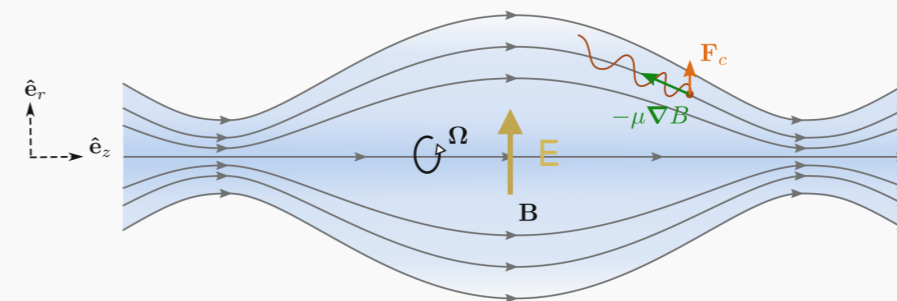


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Magnetic confinement fusion

- ▶ Centrifugal confinement in linear geometry, e. g. rotating mirrors

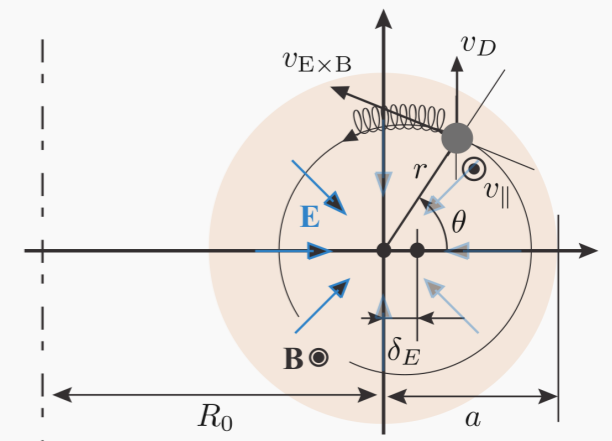


[active programs at UMD and UW]

Endrizzi et al. (2023), J. Plasma Phys. 89 (5), 975890501 - Romero-Talamas et al. (2021), APS DPP

- ▶ Magneto-electric confinement in toroidal geometry:

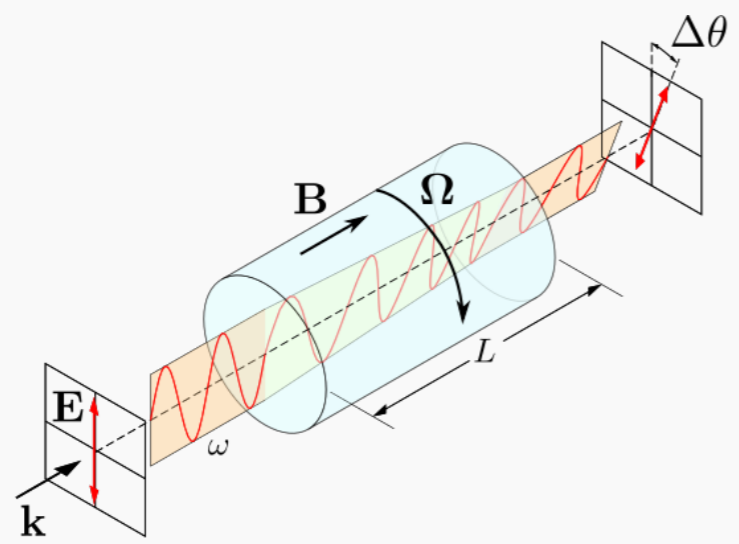
Toroidal $\mathbf{B} + \mathbf{E} \times \mathbf{B}$ poloidal rotation
 \doteq
Wave Driven Rotating Torus



Rax, Gueroult and Fisch (2017), Phys. Plasmas, **24**, 032504

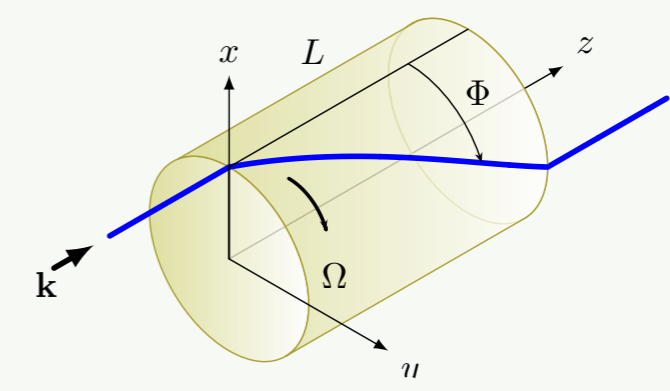
How is wave propagation affected by rotation?

Supplemental source of polarization rotation, on top of Faraday rotation



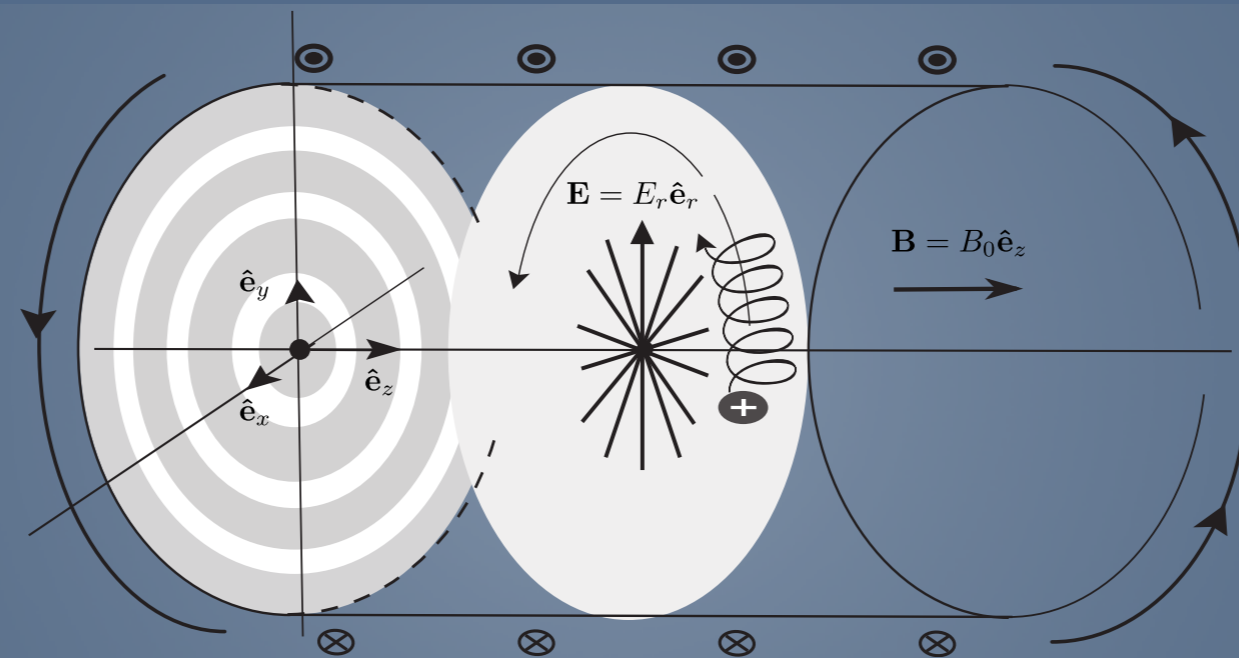
- Possibly important for pulsars, errors in galactic magnetic field estimates
Gueroult et al. (2019), Nat. Comm., **10**, 3232
- Could be used to manipulate light, tuneable non-reciprocal properties
Gueroult et al. (2020), Phys. Rev. E, **102**, 051202(R)

Light drag affects ray trajectory



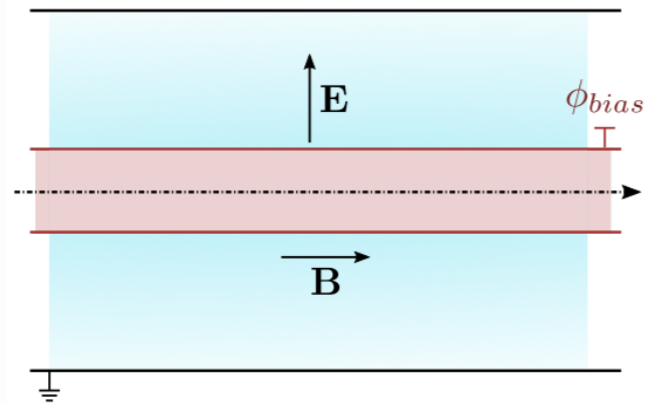
- Could have an impact on applications where accurate beam trajectory modelling is important, e. g. RF waves for MCF.
Gueroult et al. (2023), PPCF, **65**, 034006
- New diagnostics using orbital angular momentum waves
Rax and Gueroult (2021), J. Plasma Phys., **87**, 905870507

CROSS-FIELD ROTATION: HOW CAN WE CONTROL E_{\perp} ?

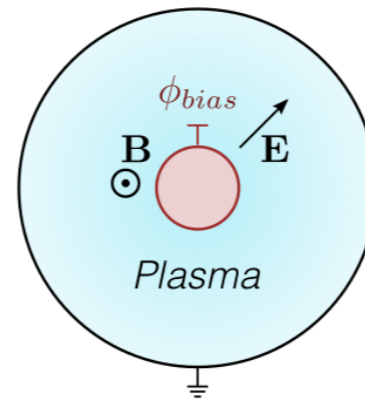


Option #1 [Simplest idea]: Central electrode

Side view



End on view

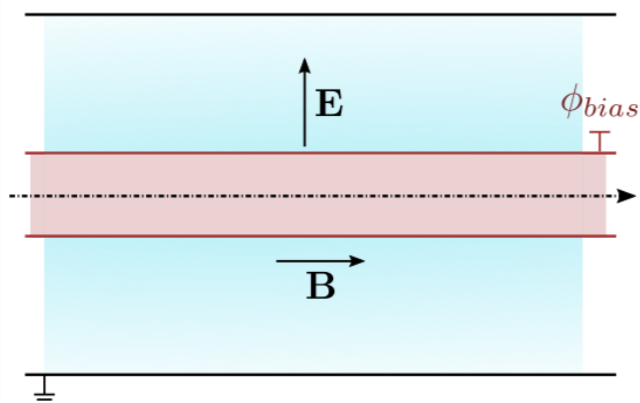


Controls $\Delta_{\perp}\phi$ but not E_{\perp}

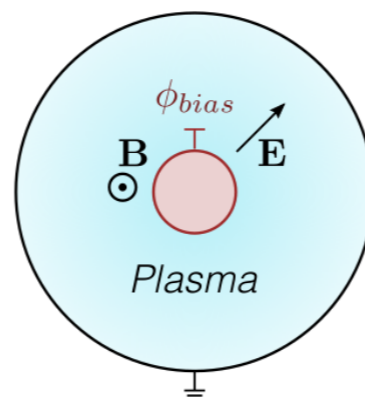
Goes against plasma's "self-organization"

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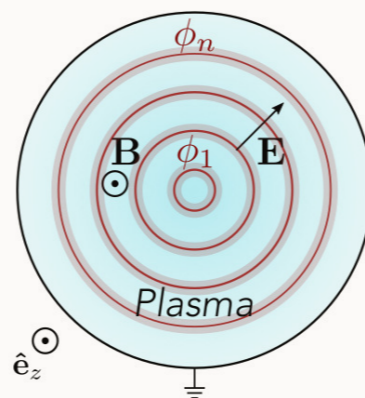
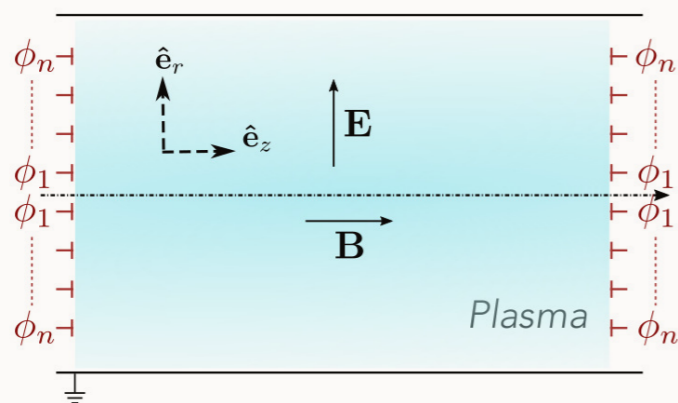
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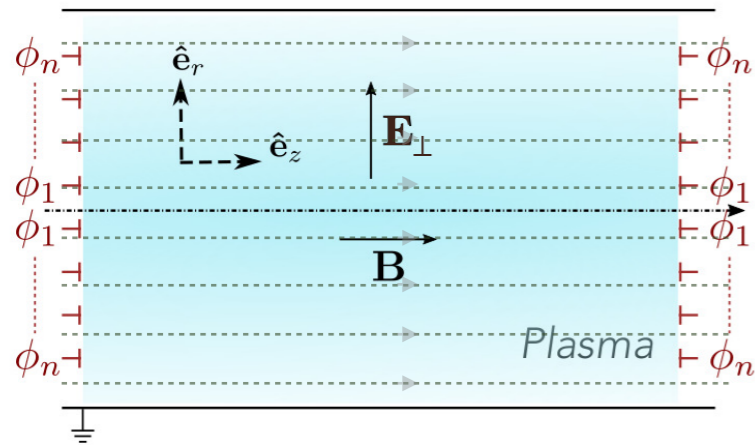
Option #2: Improved scheme



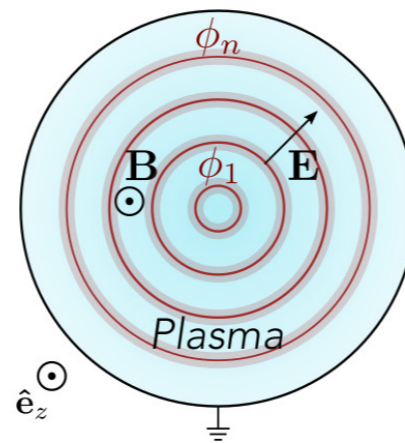
Symmetric ring-electrodes

Lehnert (1973), Phys. Scr., **7**, 102

Side view



End on view



Experiment

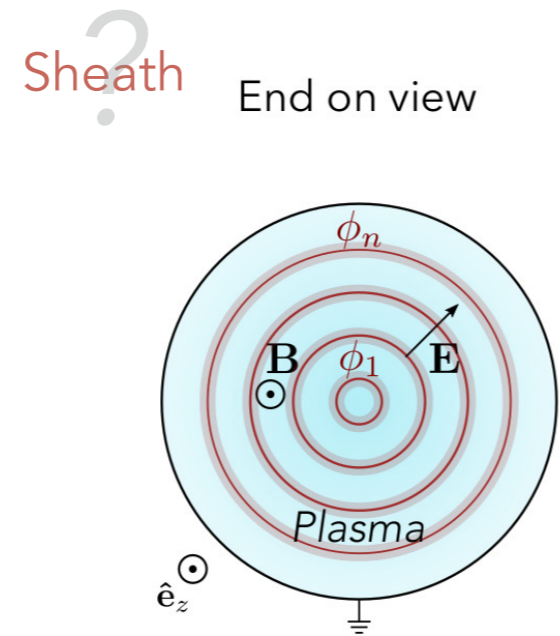
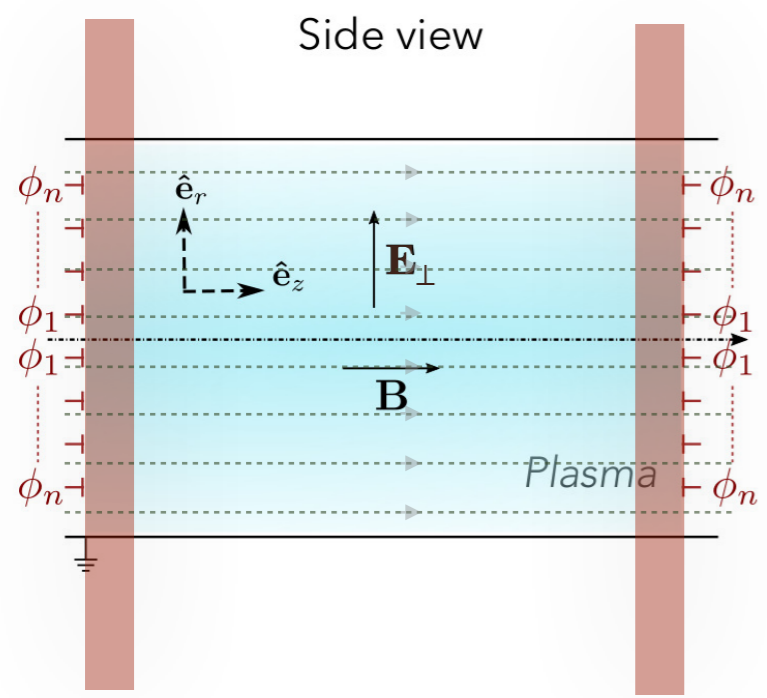
CONTROL

- Q-machine, West Virginia Univ.
- Large diameter helicon, Kyushu Univ.
- QT-Upgrade machine, Tohoku Univ.
- Gamma 10 mirror, Univ. Tsukuba
- LAPD afterglow, Univ. California Los Angeles
- Phaedrus tandem mirror, Univ. Wisconsin
- Helcat, Univ. New Mexico
- PMFX, Princeton Plasma Physics Lab.
- C-2 device, Tri Alpha Energy
- KMAX mirror, Univ. Sci. Tech. China



Experimental results so far are however inconclusive: *when can end-electrodes biasing be expected to be effective?*

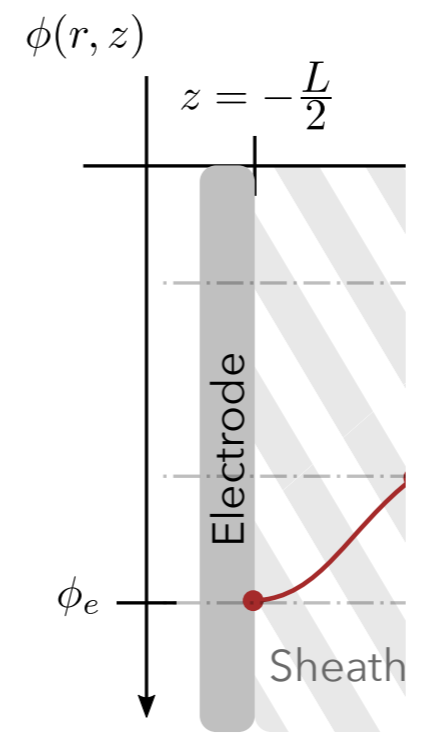
Gueroult et al. (2019), Phys. Plasmas, **26**, 122106

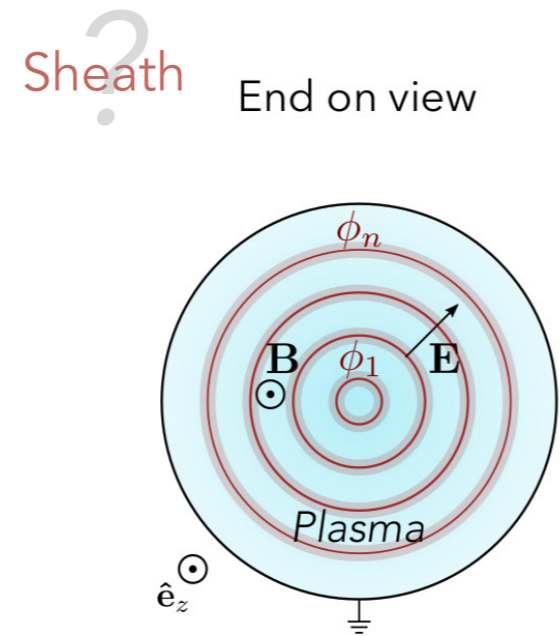
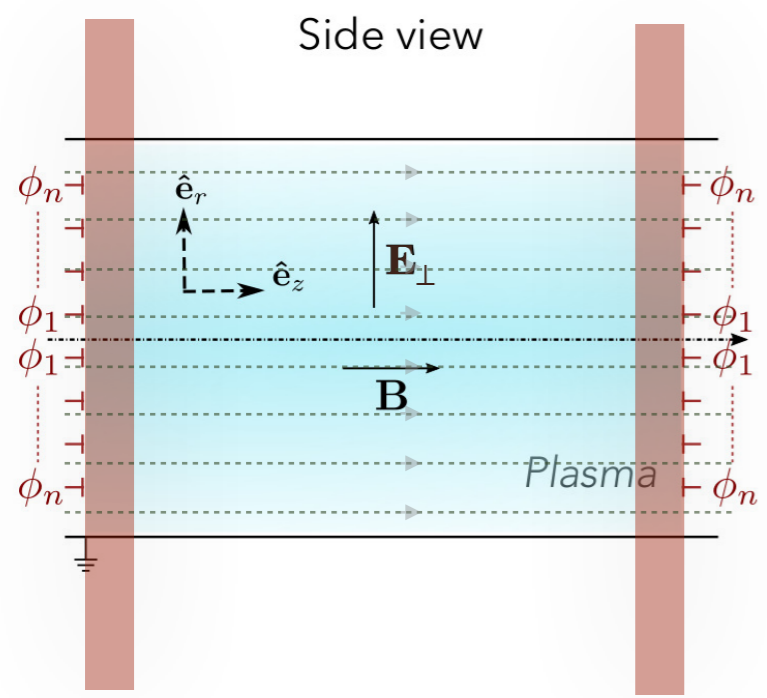


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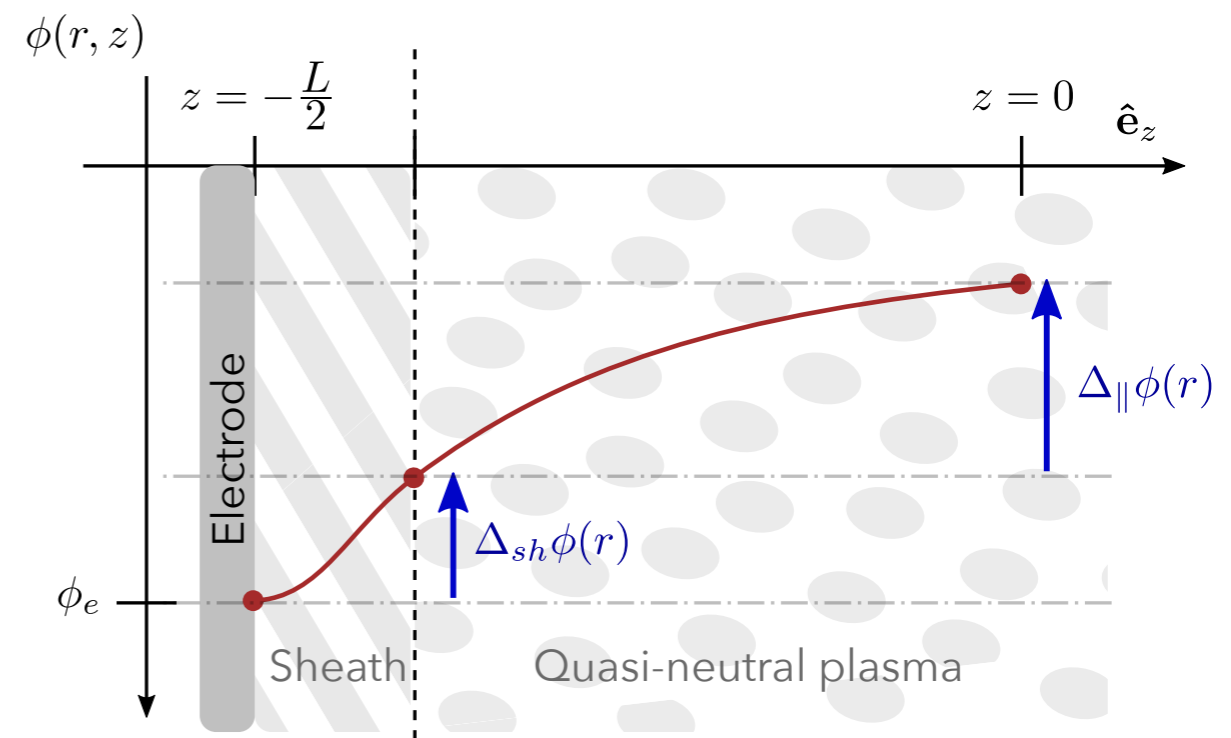


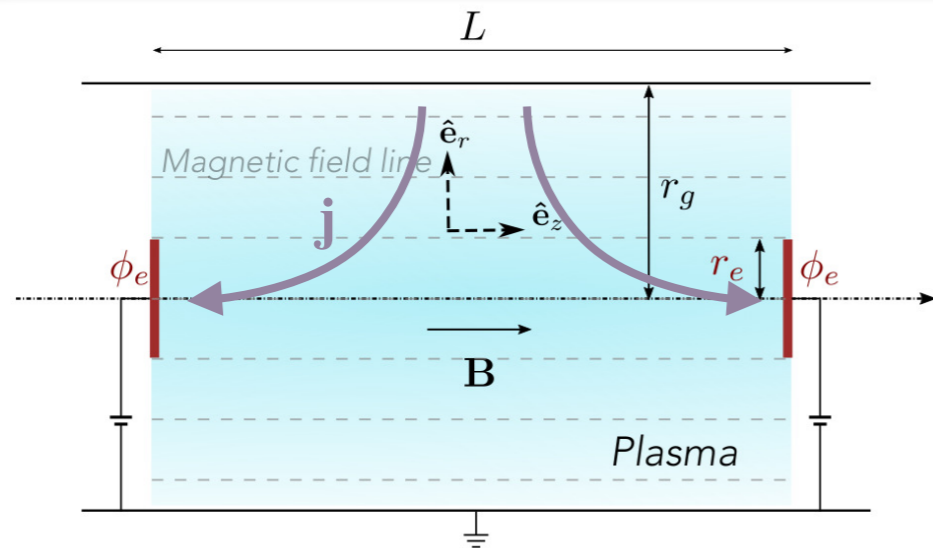
B line =
isopotential

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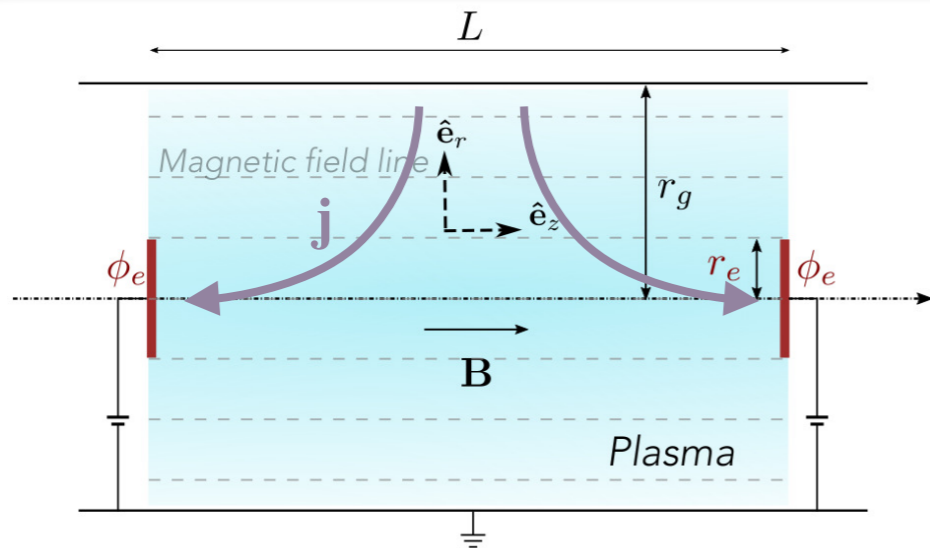
(a) Symmetric button end-electrodes

Assumptions:

- ▶ Negative bias, ion sheath, $\sigma_{\parallel} \rightarrow \infty$
- ▶ Classical transport, radial resistance

$$R_p \doteq \frac{1}{2\pi L \sigma_{\perp}} \ln \left(\frac{r_g}{r_e} \right)$$

Liziakin *et al.*, PSST (2020), 29, 015008



(a) Symmetric button end-electrodes

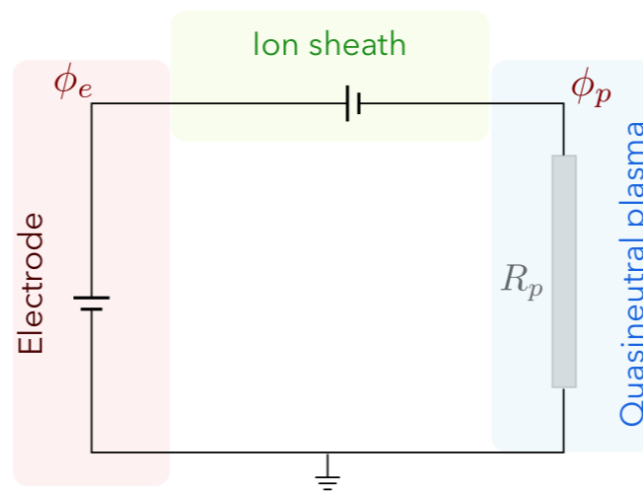
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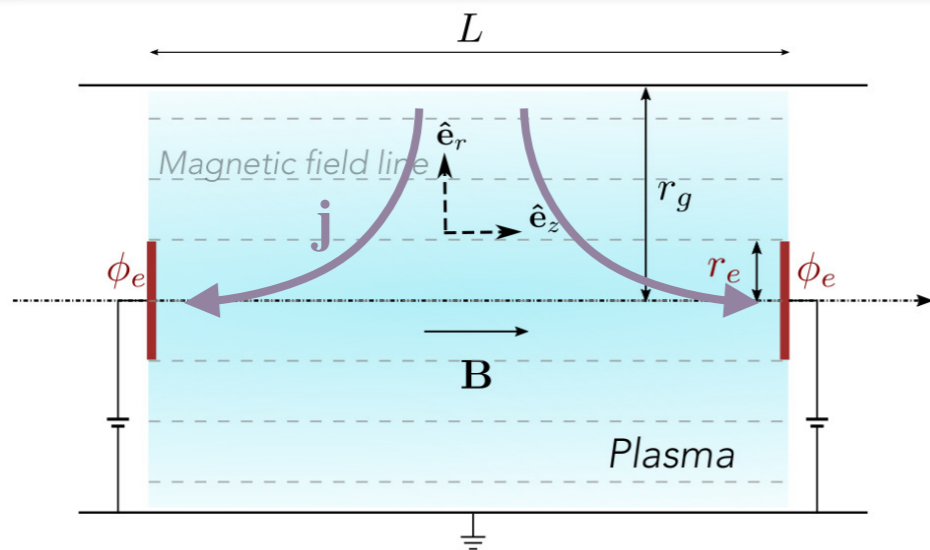
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Equivalent circuit:





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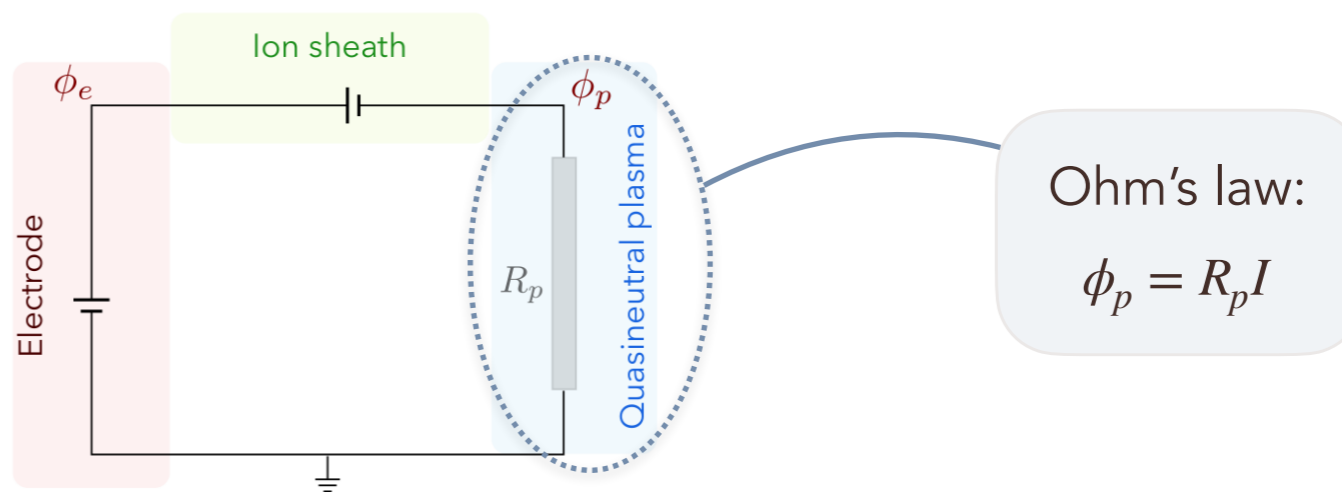
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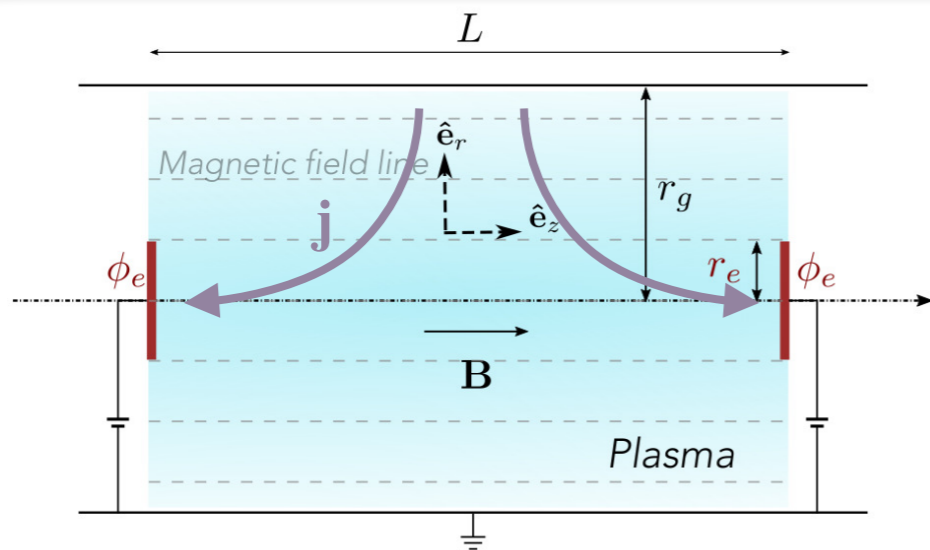
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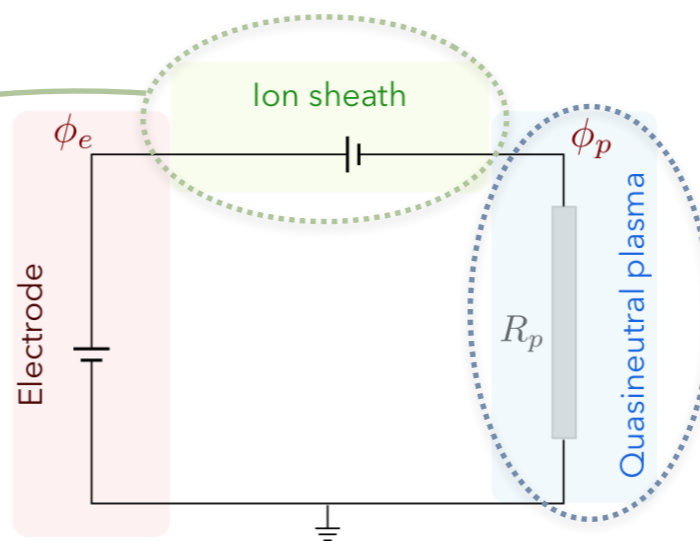
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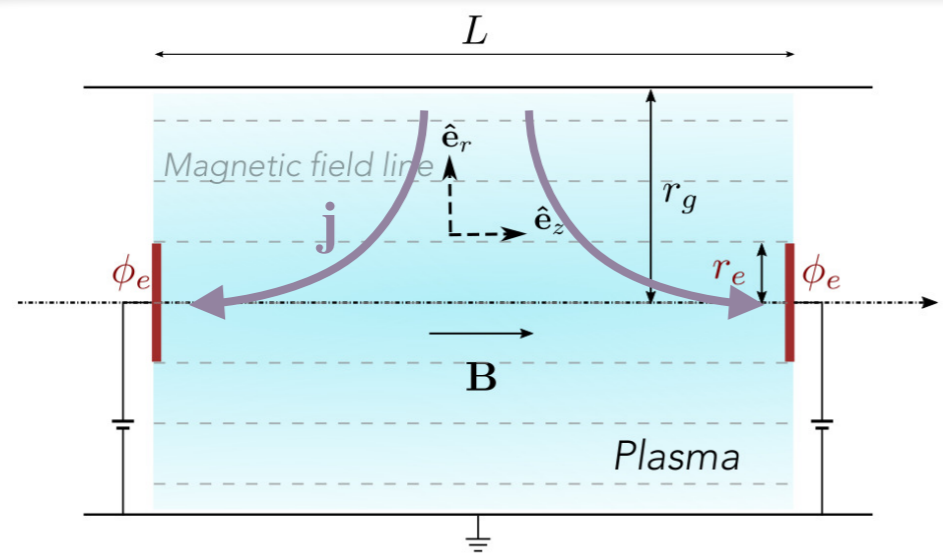
Current balance at electrode:

$$I = I_{is} \left[1 - \exp \left(\Lambda + e \frac{\phi_e - \phi_p}{T_e} \right) \right]$$



Ohm's law:

$$\phi_p = R_p I$$



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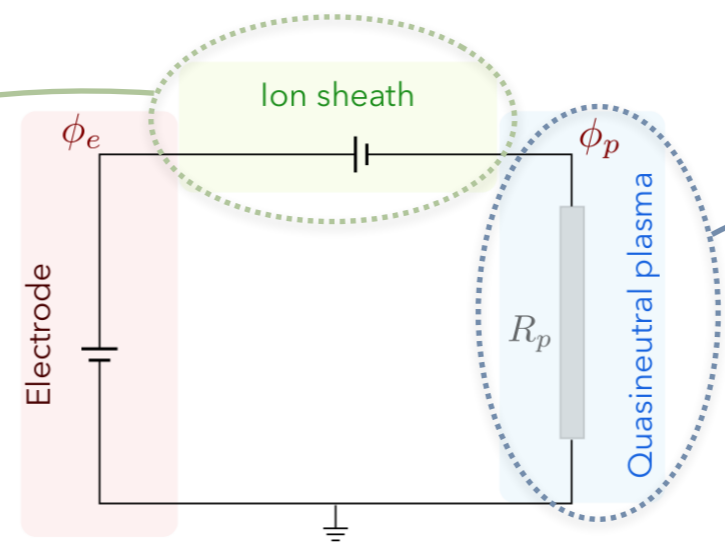
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Equating currents gives a transcendental equation for the plasma potential

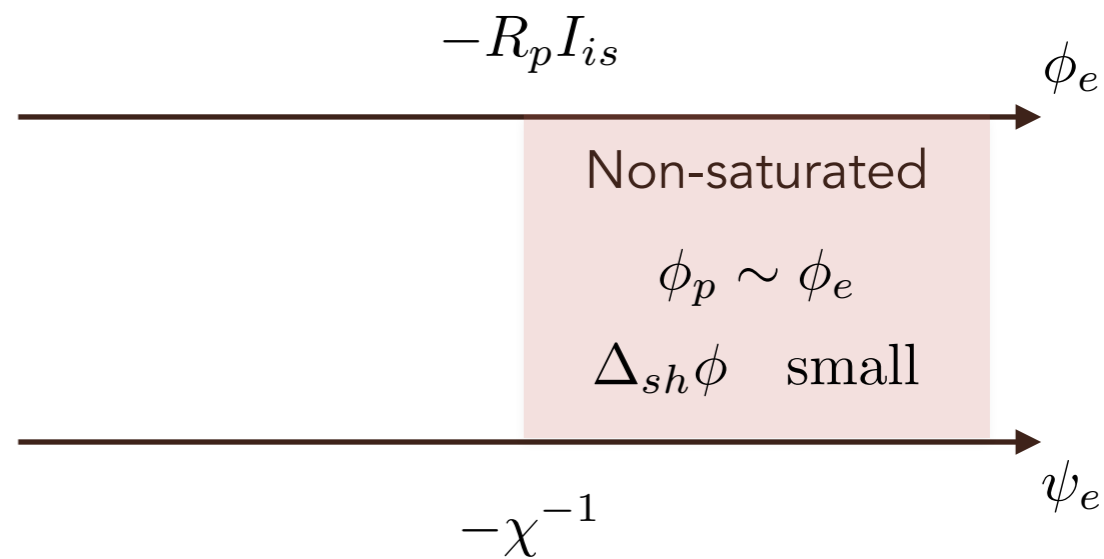
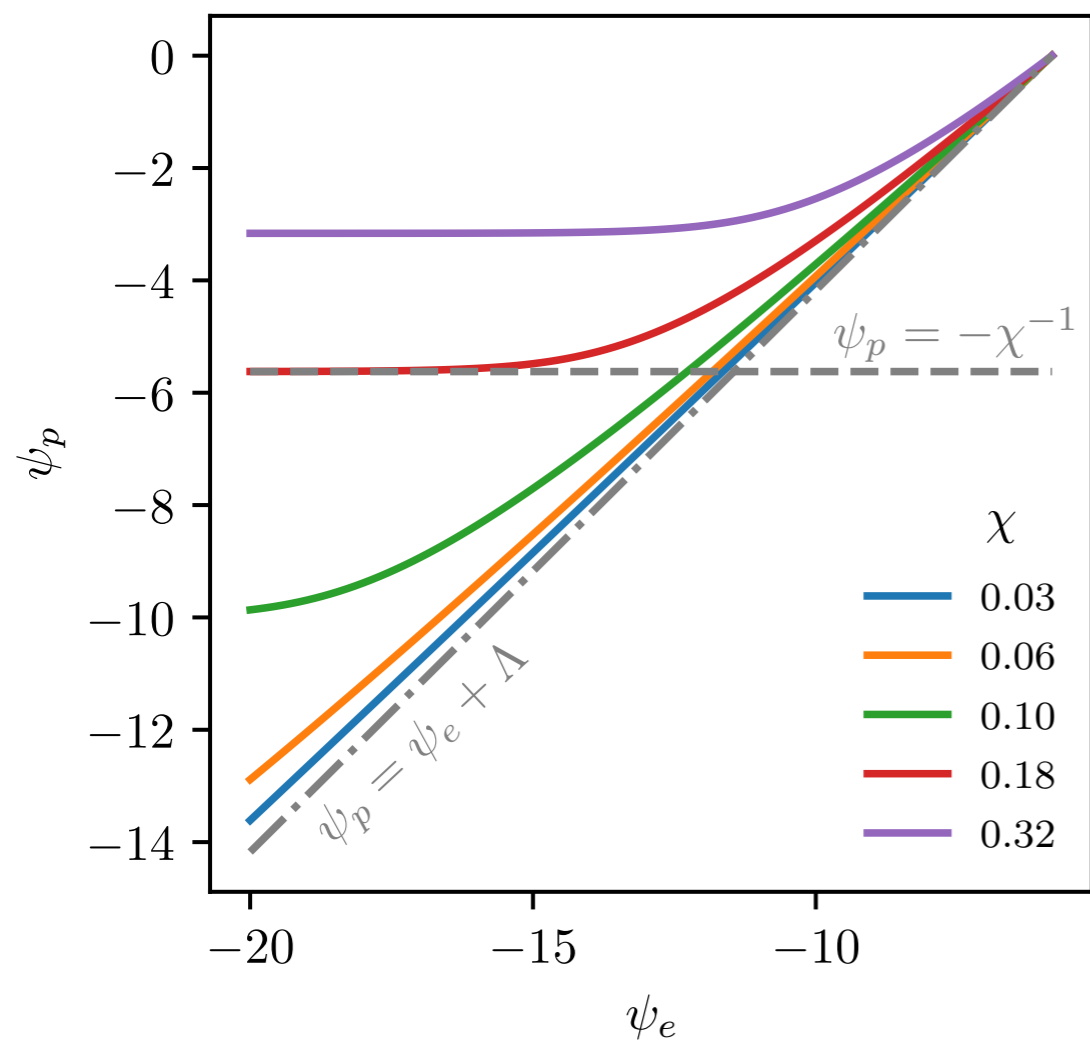
$$\exp(\Lambda + \psi_e - \psi_p) - 1 - \chi \psi_p = 0$$

$$\left| \begin{array}{l} \psi \doteq e\phi/T_e \\ \chi \doteq T_e/(eRI_{is}) \end{array} \right.$$

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$$\psi \doteq eV/T_e$$

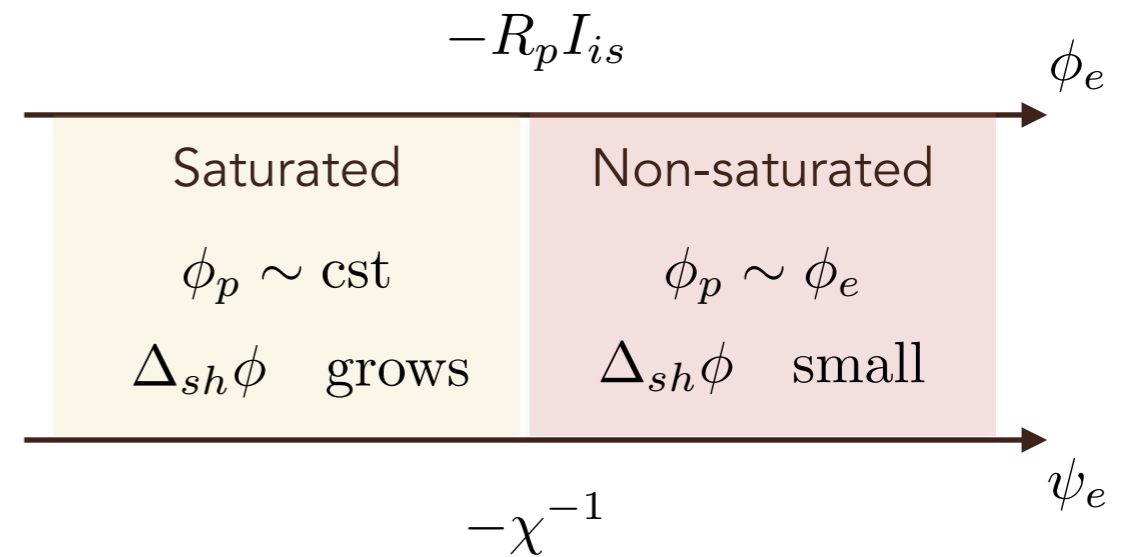
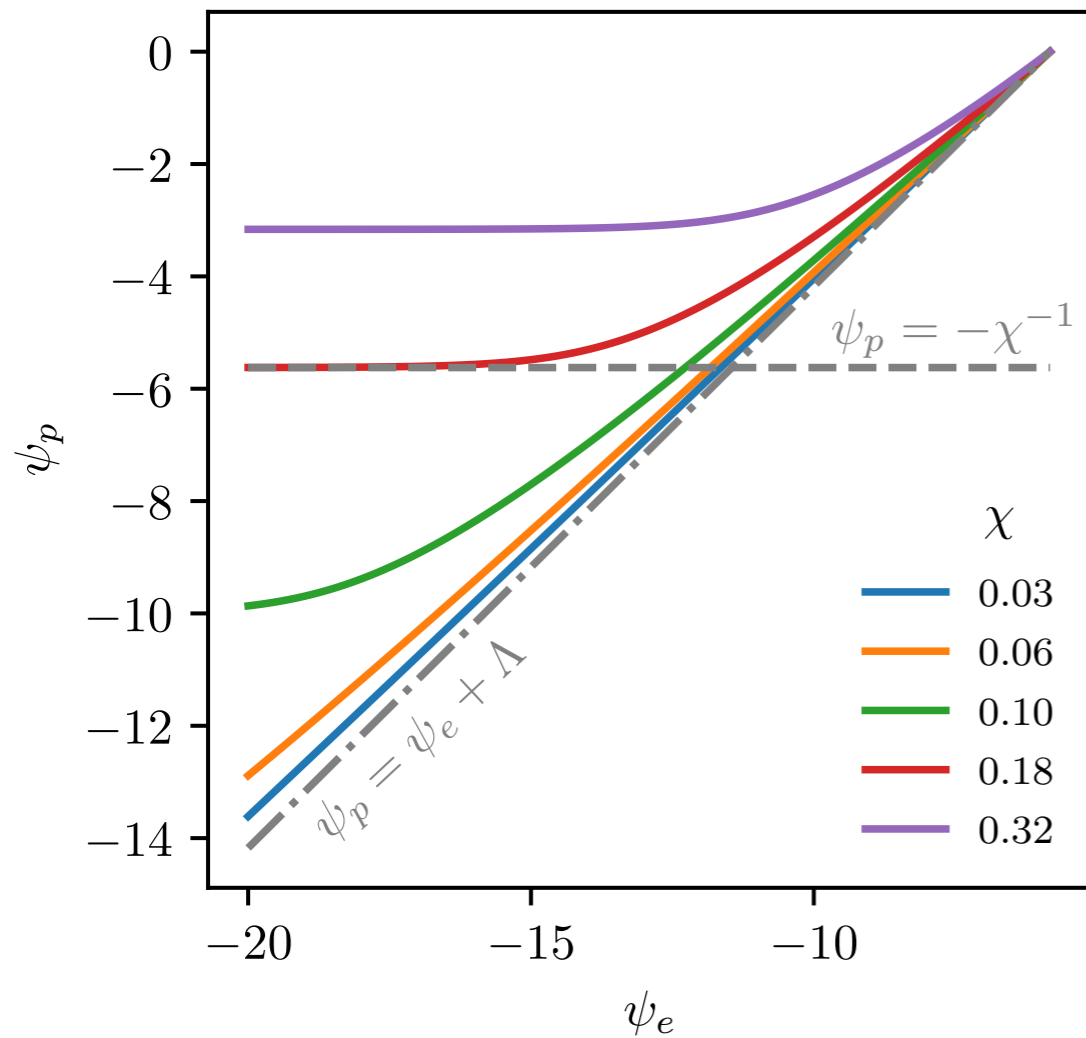
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Trotabas and Gueroult PSST (2022), 31, 025001

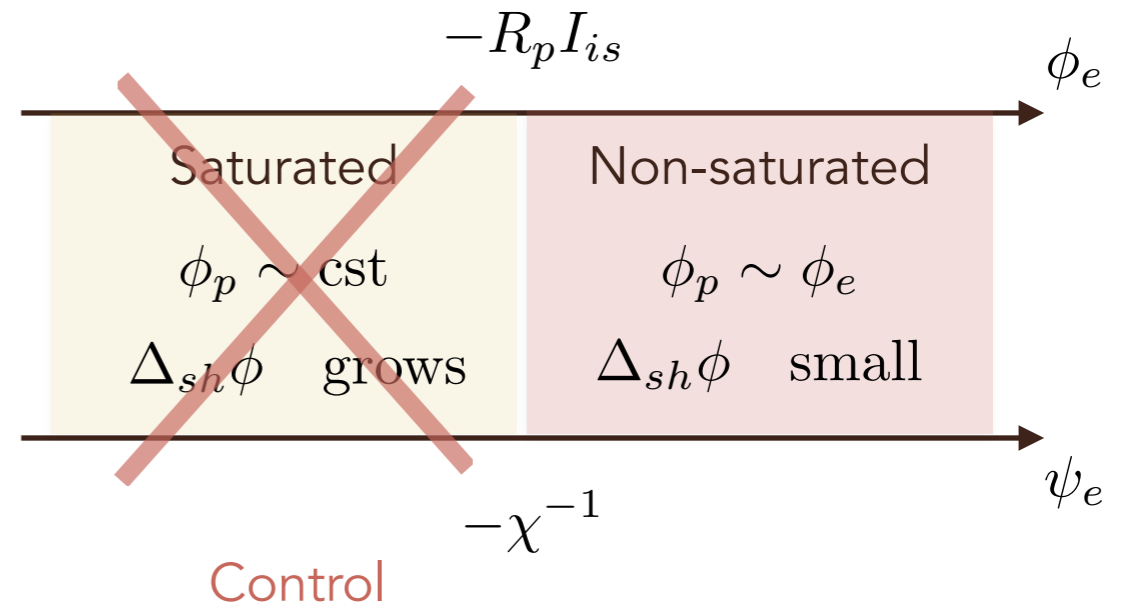
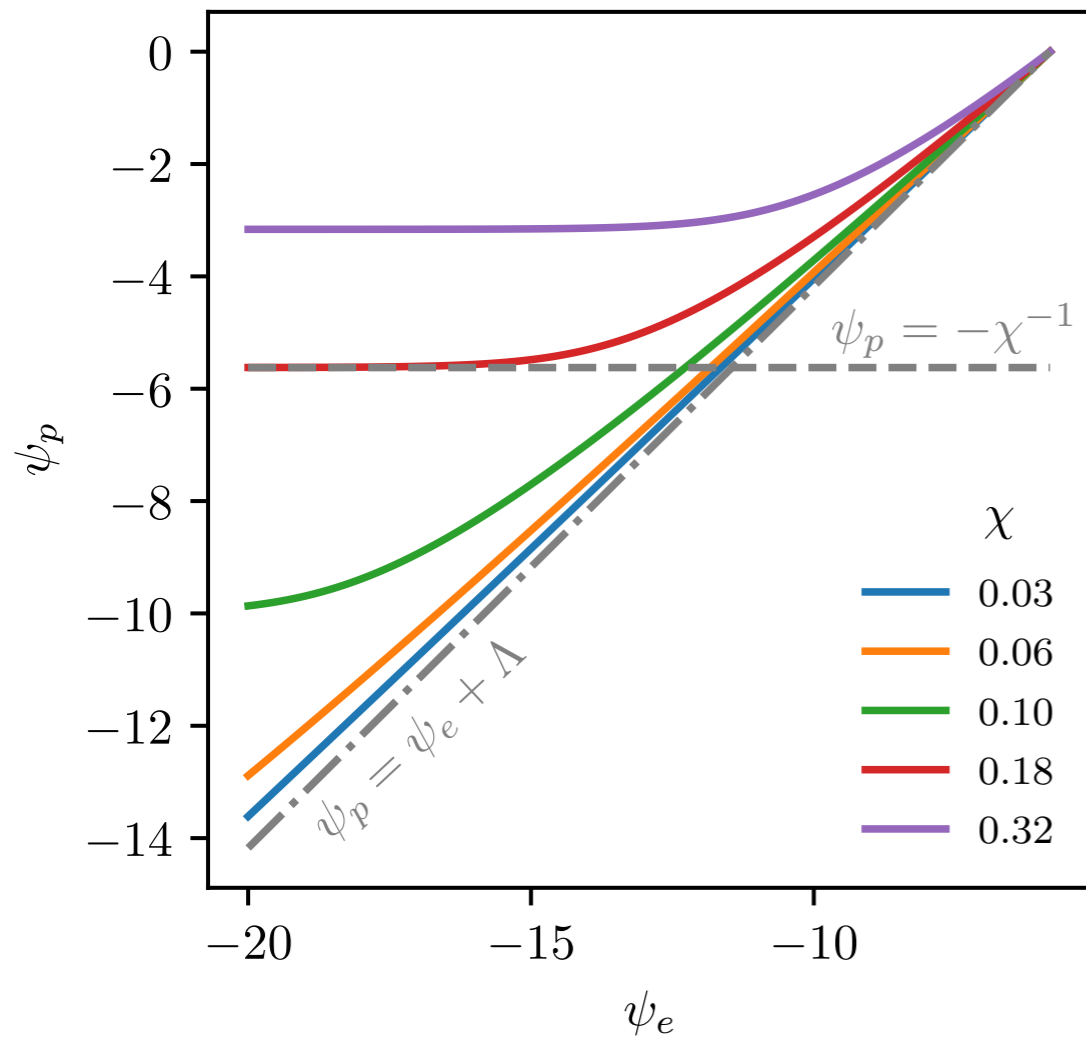
Trotabas, PhD Thesis UT3 Paul Sabatier (2022)

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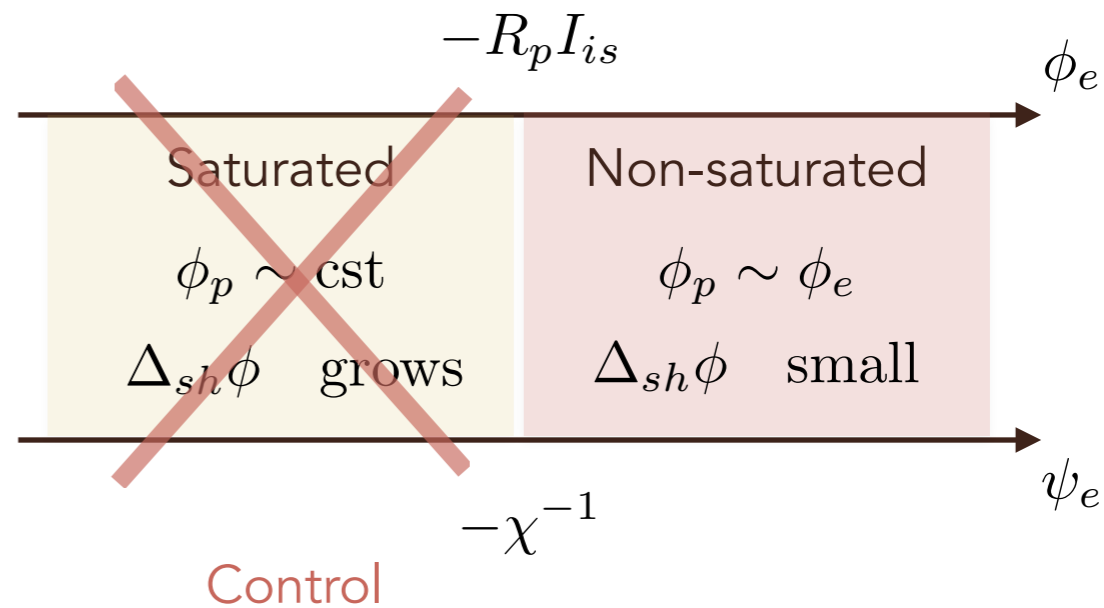
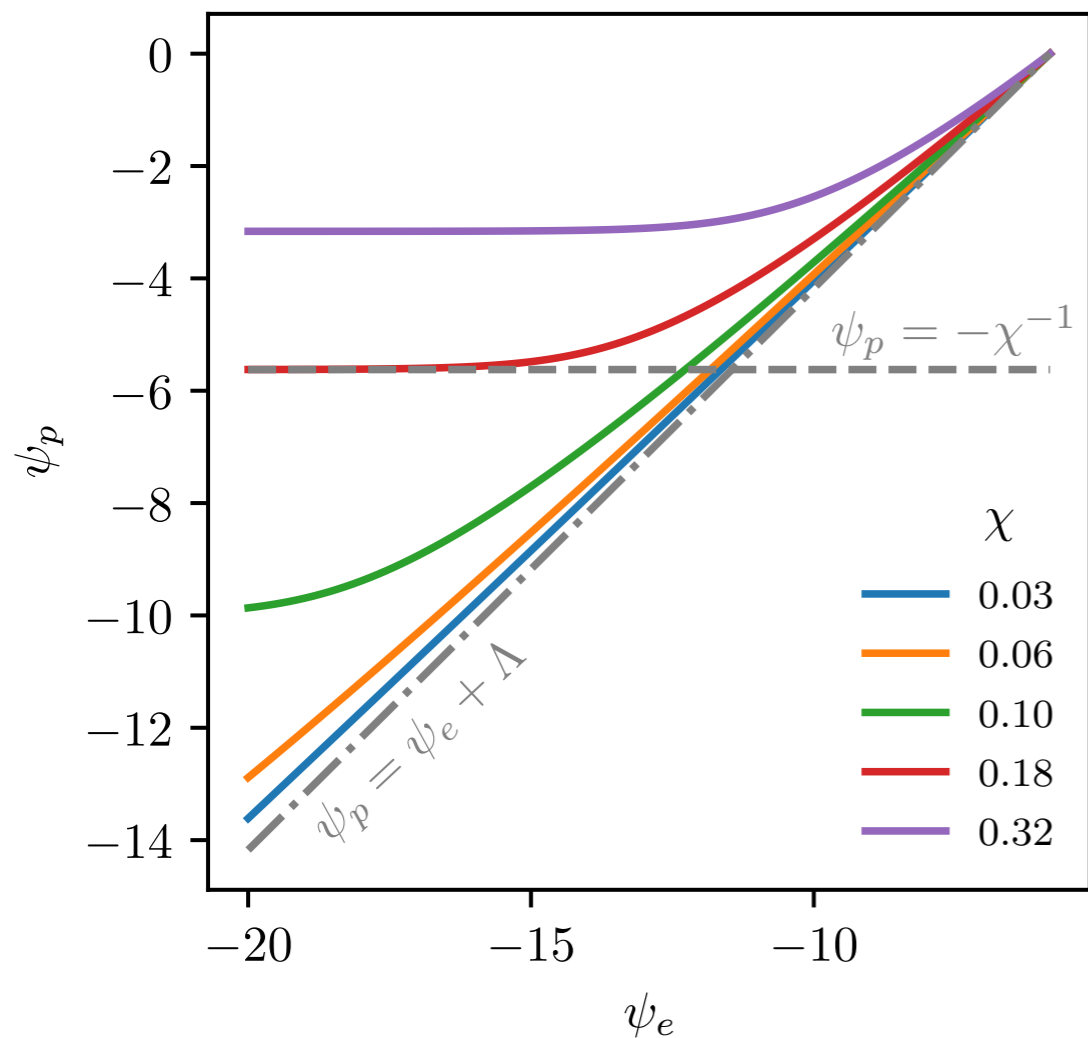
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Trotabas and Gueroult PSST (2022), 31, 025001

Trotabas, PhD Thesis UT3 Paul Sabatier (2022)

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Model can straightforwardly be extended to a heated cathode (thermionic emission)

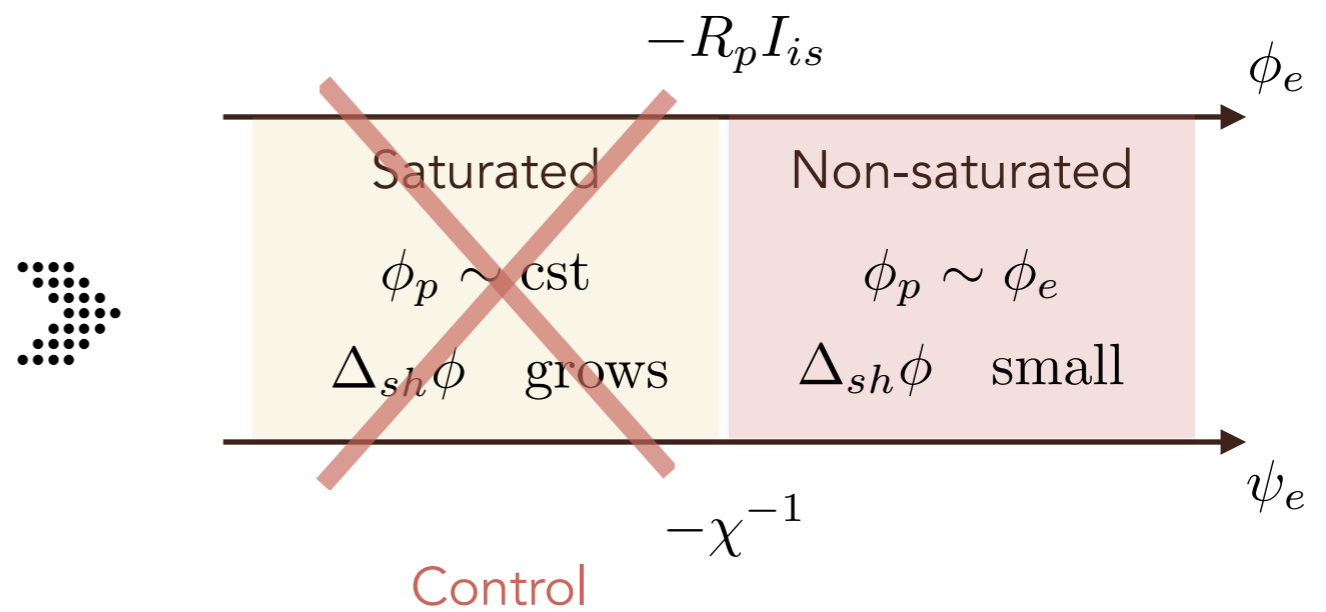
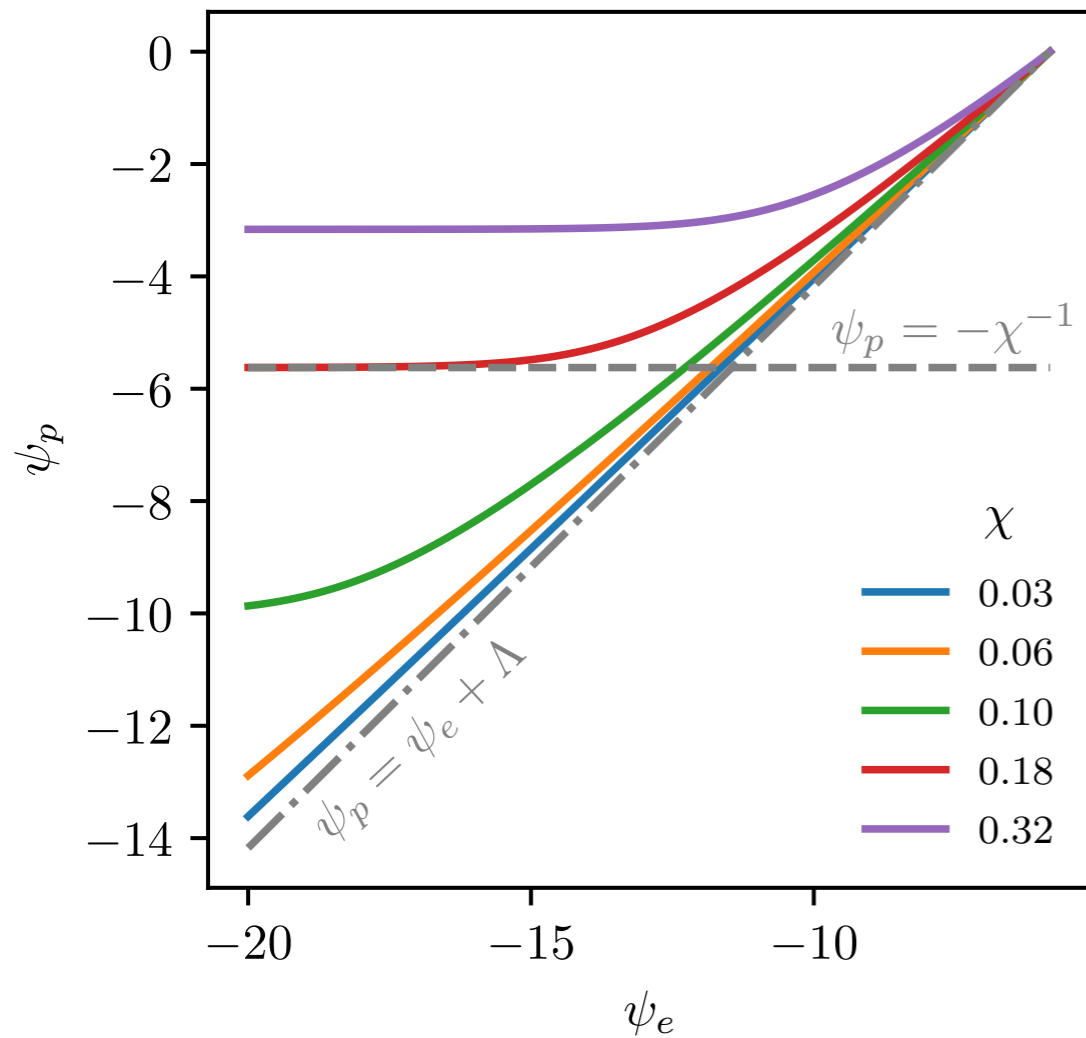
$$\exp(\Lambda + \psi_e - \psi_p) - 1 - \Xi - \chi\psi_p = 0$$

\uparrow
 $\frac{j_{eth}}{j_{is}}$

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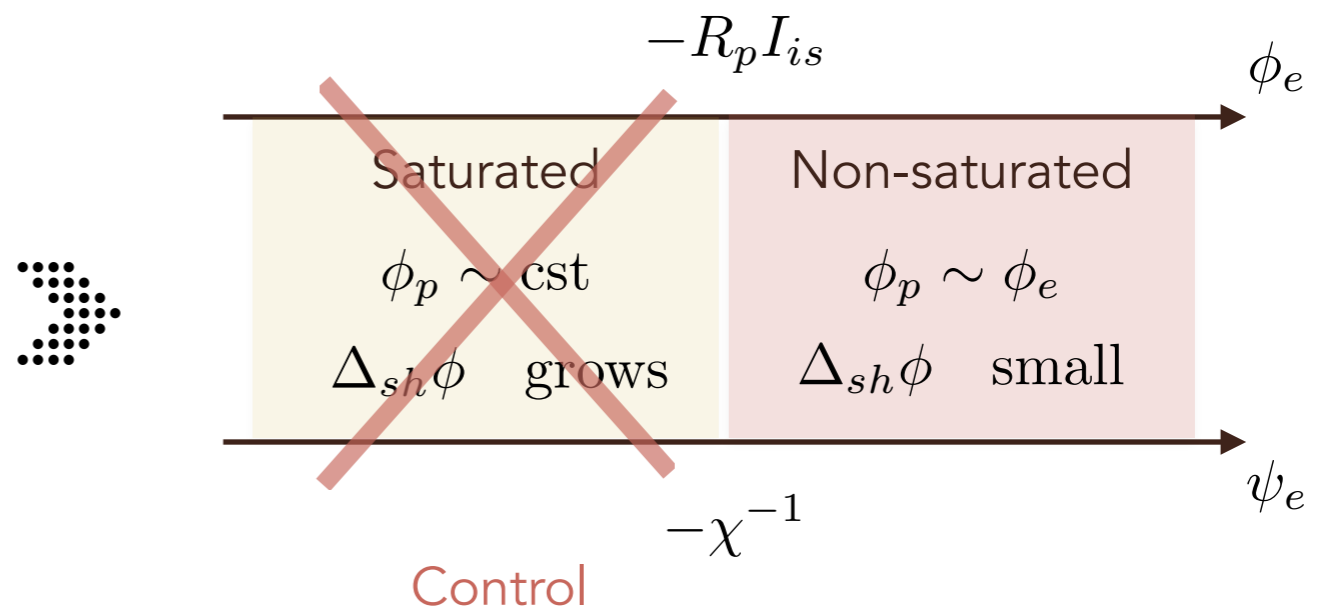
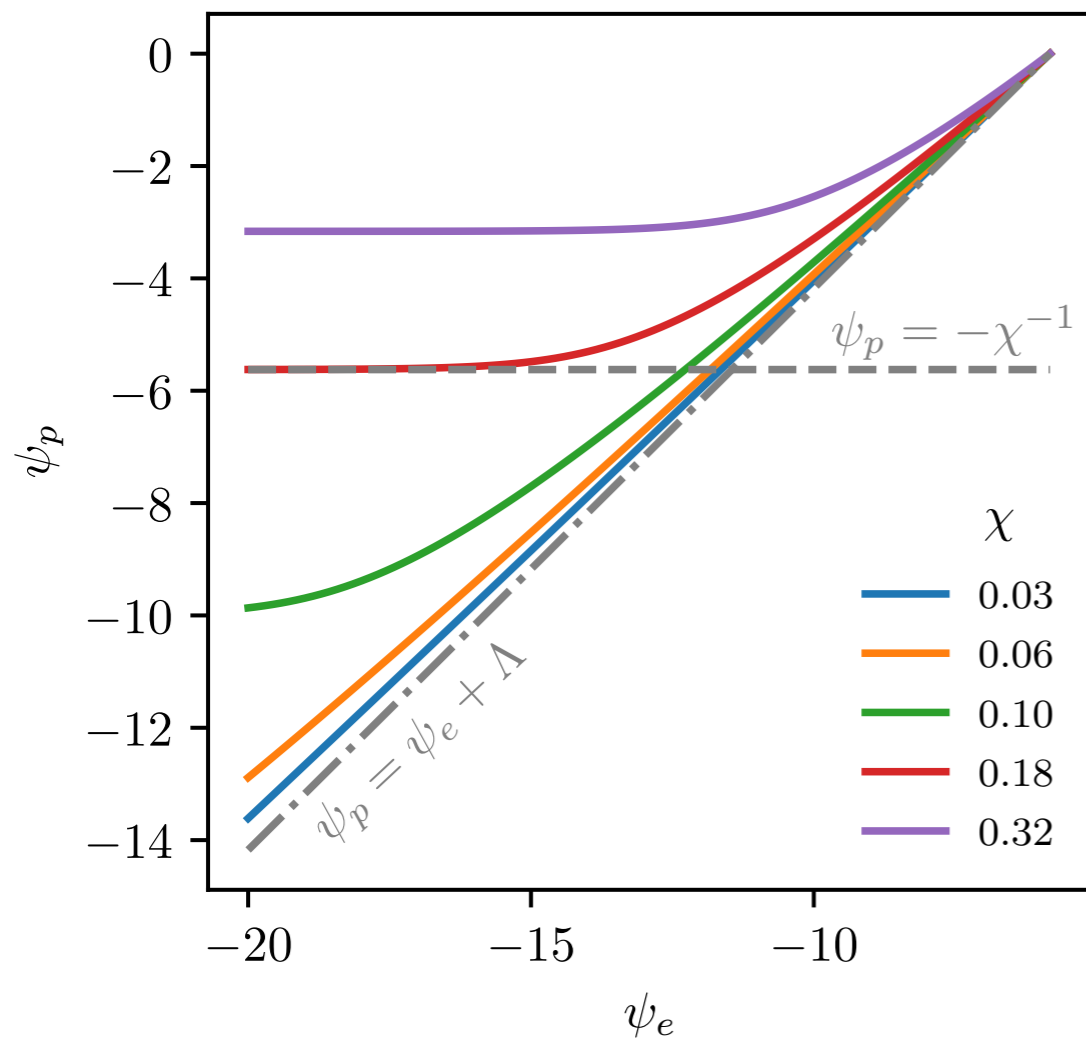


$$\max |\psi_p| = \frac{1 + \Xi}{\chi} = (1 + \Xi) \max |\psi_p^{cold}|$$

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Trotabas and Gueroult PSST (2022), 31, 025001

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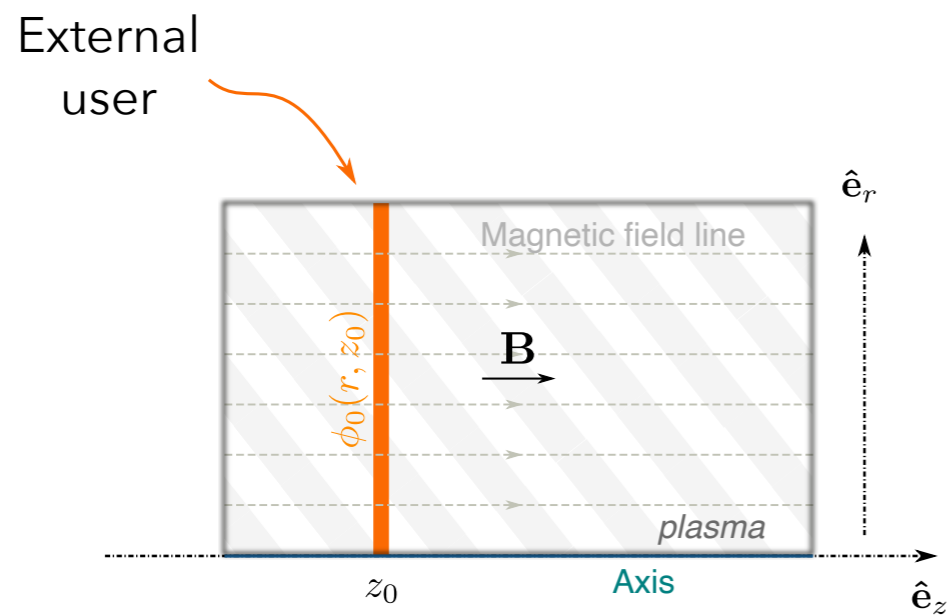
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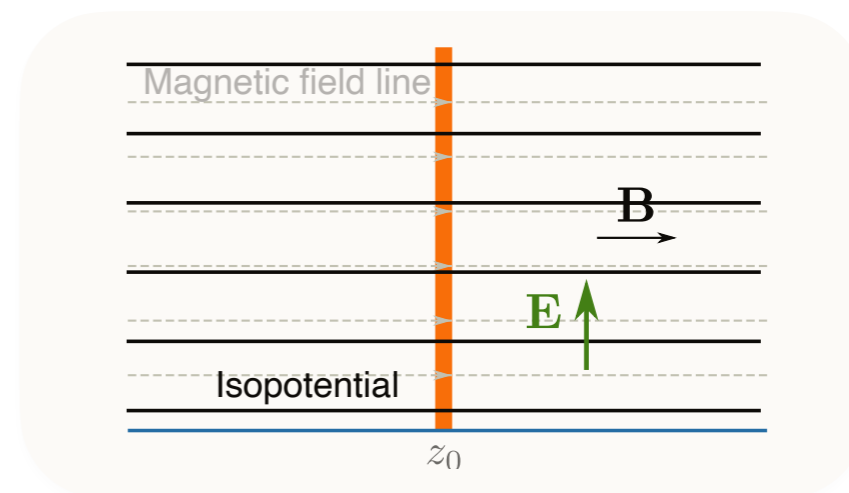
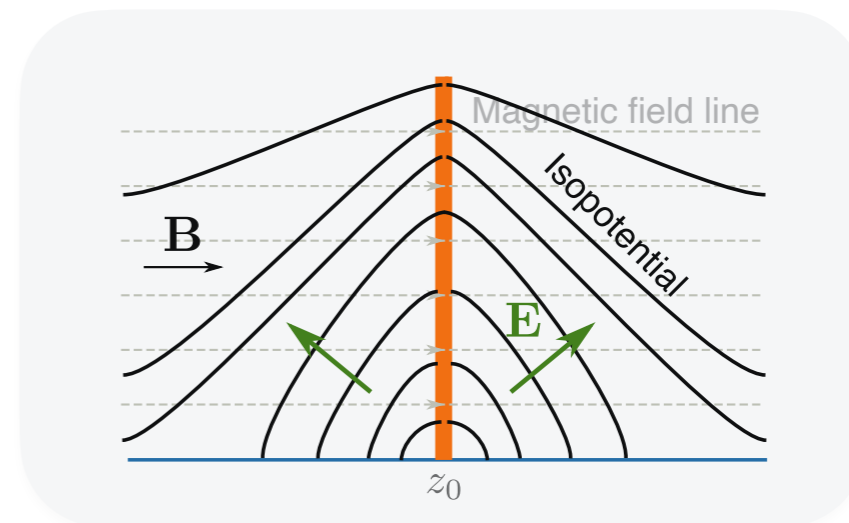
Improved control, plasma potential can be \searrow by a factor j_{eth}/j_{is}
 Consistent with experiments. *More on this in N. Plihon's talk.*

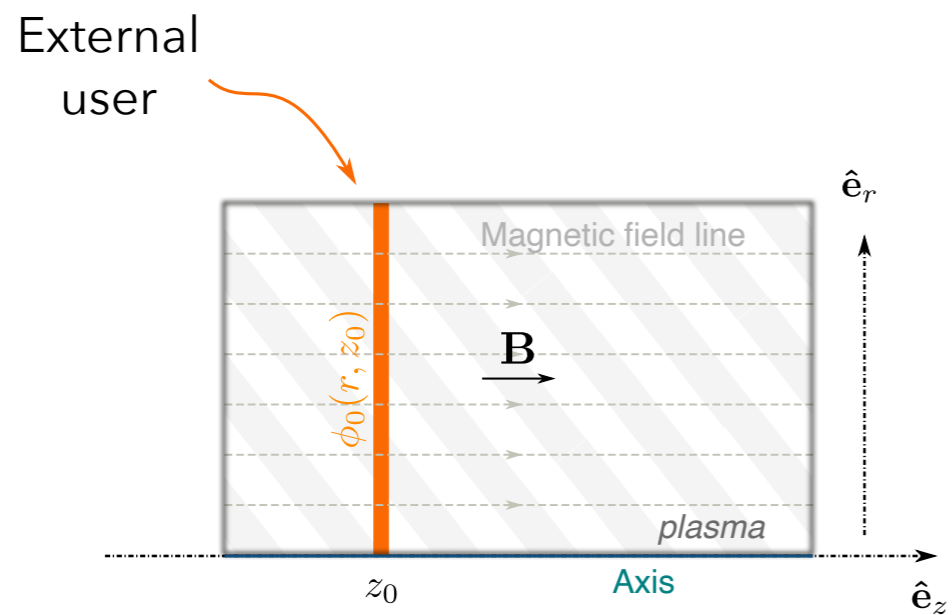


vacuum field

?

infinite conductivity along field lines ($\sigma_{\parallel}/\sigma_{\perp} \rightarrow \infty$)

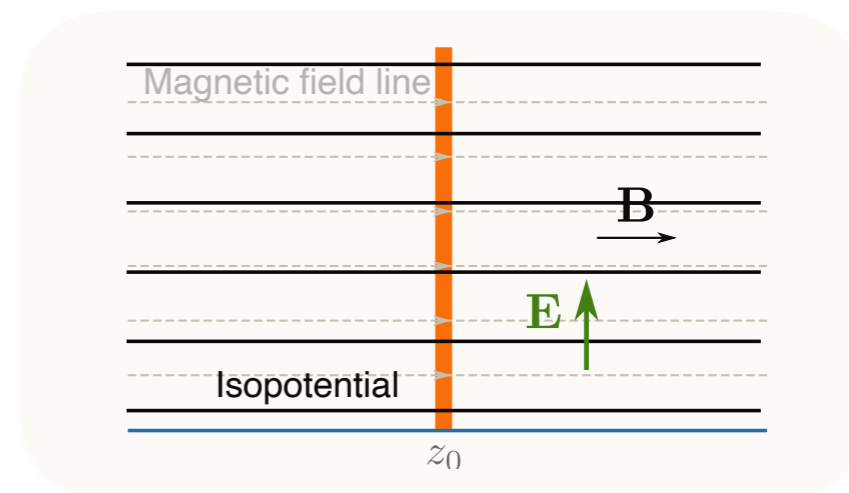
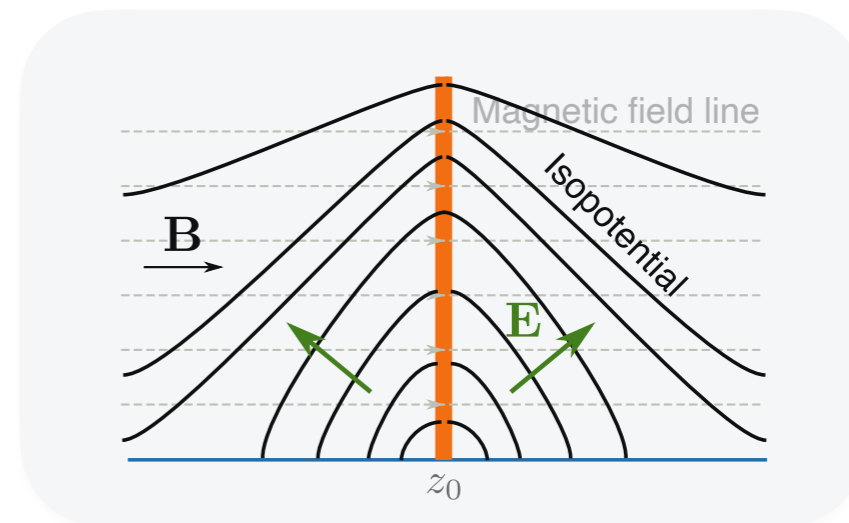




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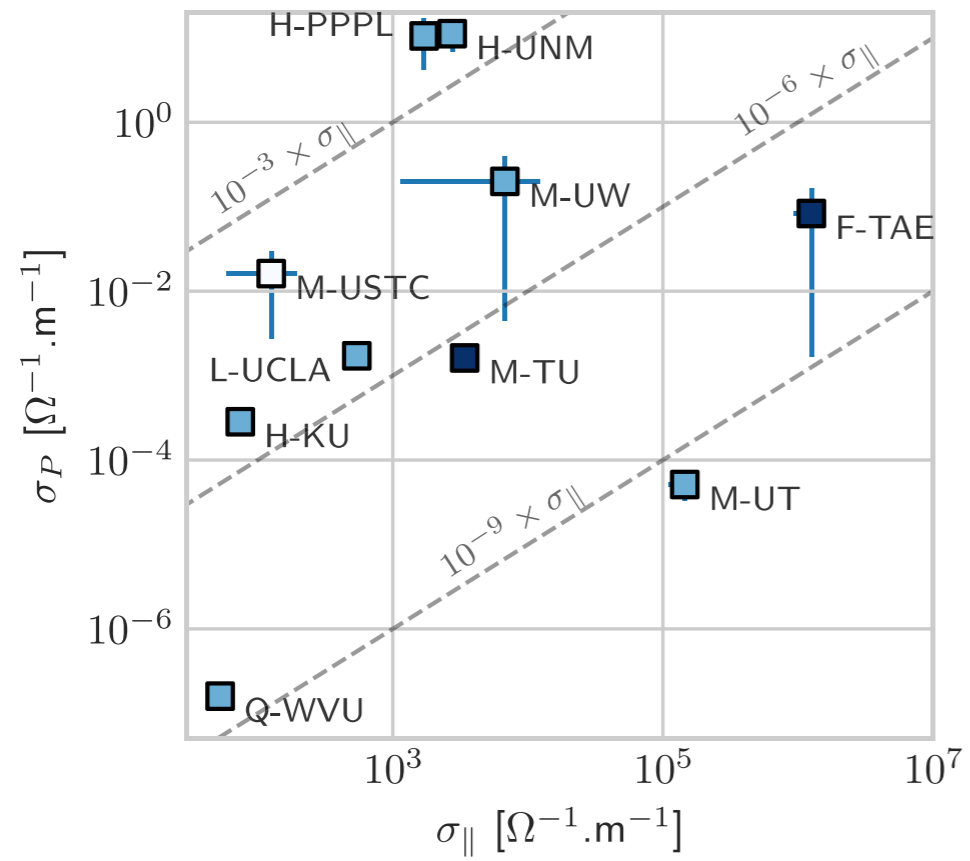
Do we really need $\sigma_{\parallel} \rightarrow \infty$? Probably not

$$\tau = \sqrt{\frac{\sigma_{\perp}}{\sigma_{\parallel}}} \frac{L}{a} \ll 1$$

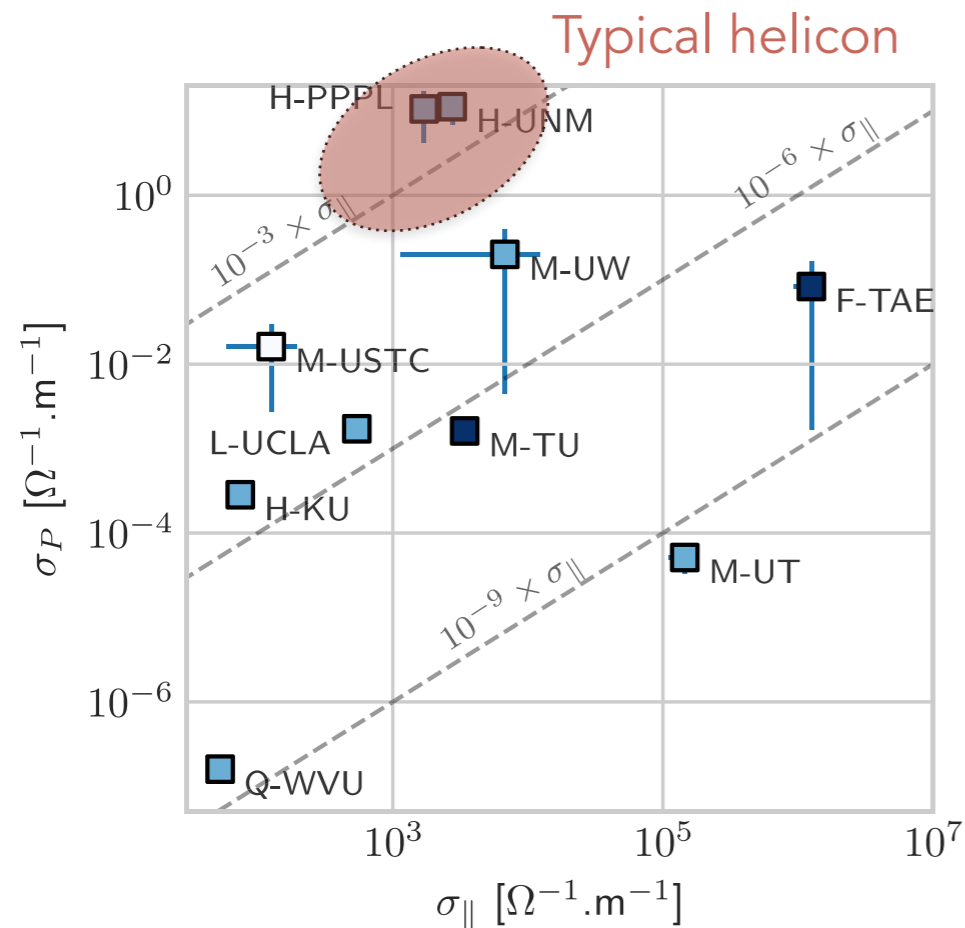
Length

Radius

Poulos (2019), Phys. Plasmas, **26**, 22104
 Trotabas and Gueroult (2022), PSST, **31**, 025001

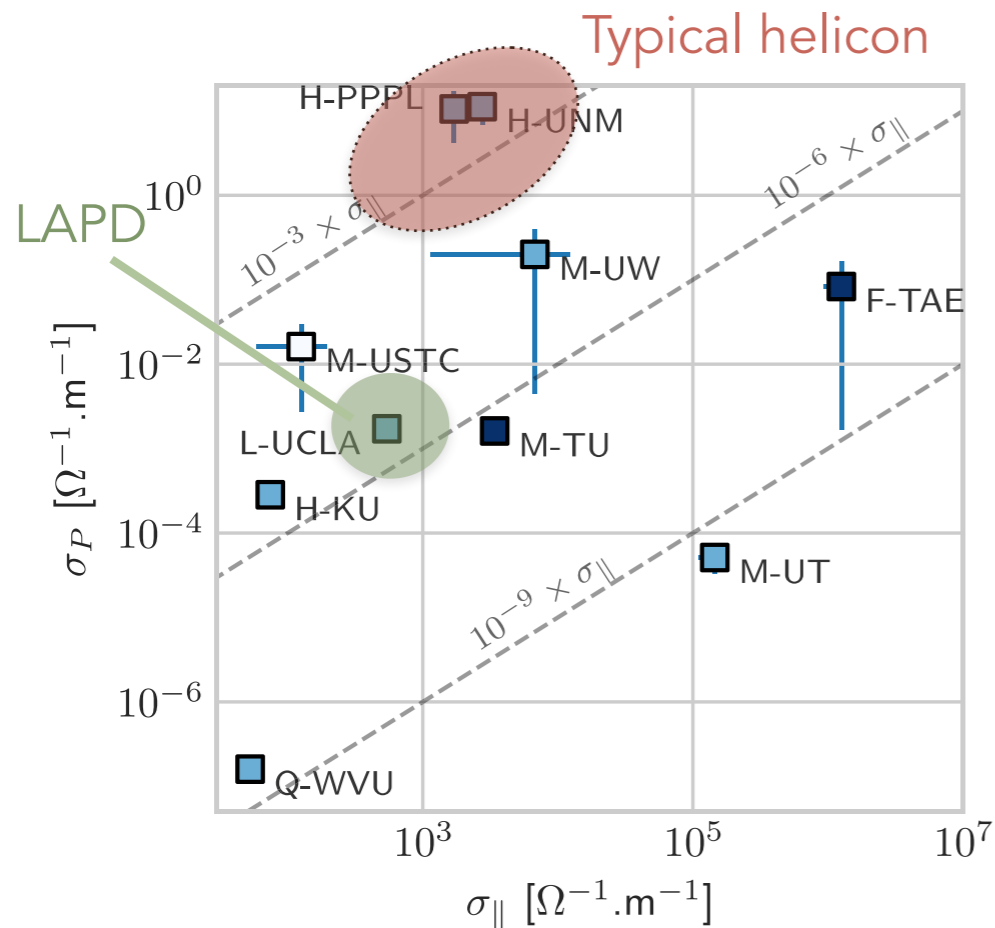


Gueroult et al. (2016), PSST, 25, 035024



- Helicons tend to have not so small $\sigma_{\perp}/\sigma_{\parallel}$. Results from comparatively larger neutral fill pressure ($p \sim$ few mTorr).
- Other sources lead to smaller $\sigma_{\perp}/\sigma_{\parallel}$, and thus lower axial voltage drop along field lines.
- Example is hot cathode on LAPD ($p < 3.5 \cdot 10^{-2}$ mTorr).

Gueroult et al. (2016), PSST, 25, 035024



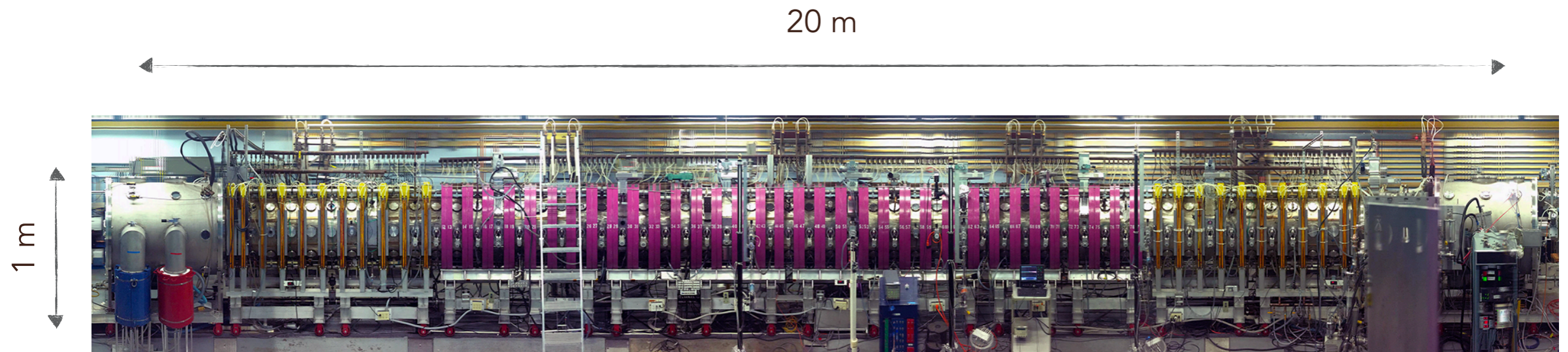
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Gueroult et al. (2016), PSST, 25, 035024

EXPERIMENT ON THE LARGE PLASMA DEVICE (LAPD) AT UCLA

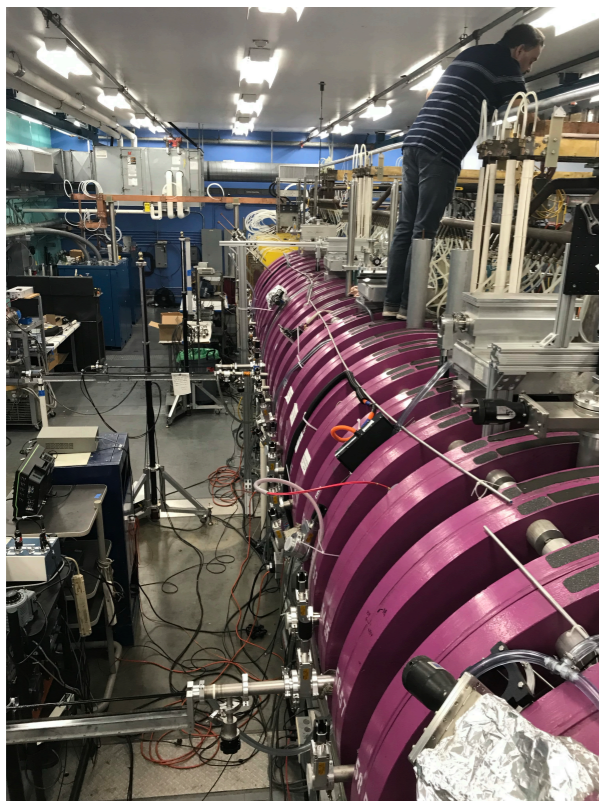
IN COLLABORATION WITH
S. TRIPATHI, N. J. FISCH

AND WITH HELP FROM UCLA BAPSF GROUP

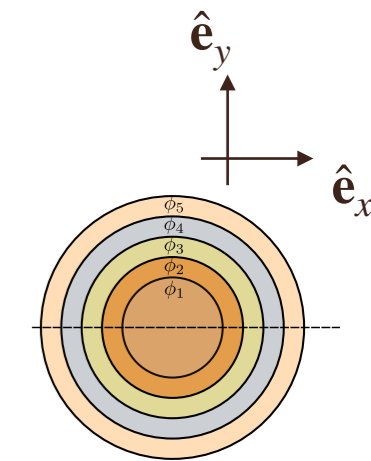
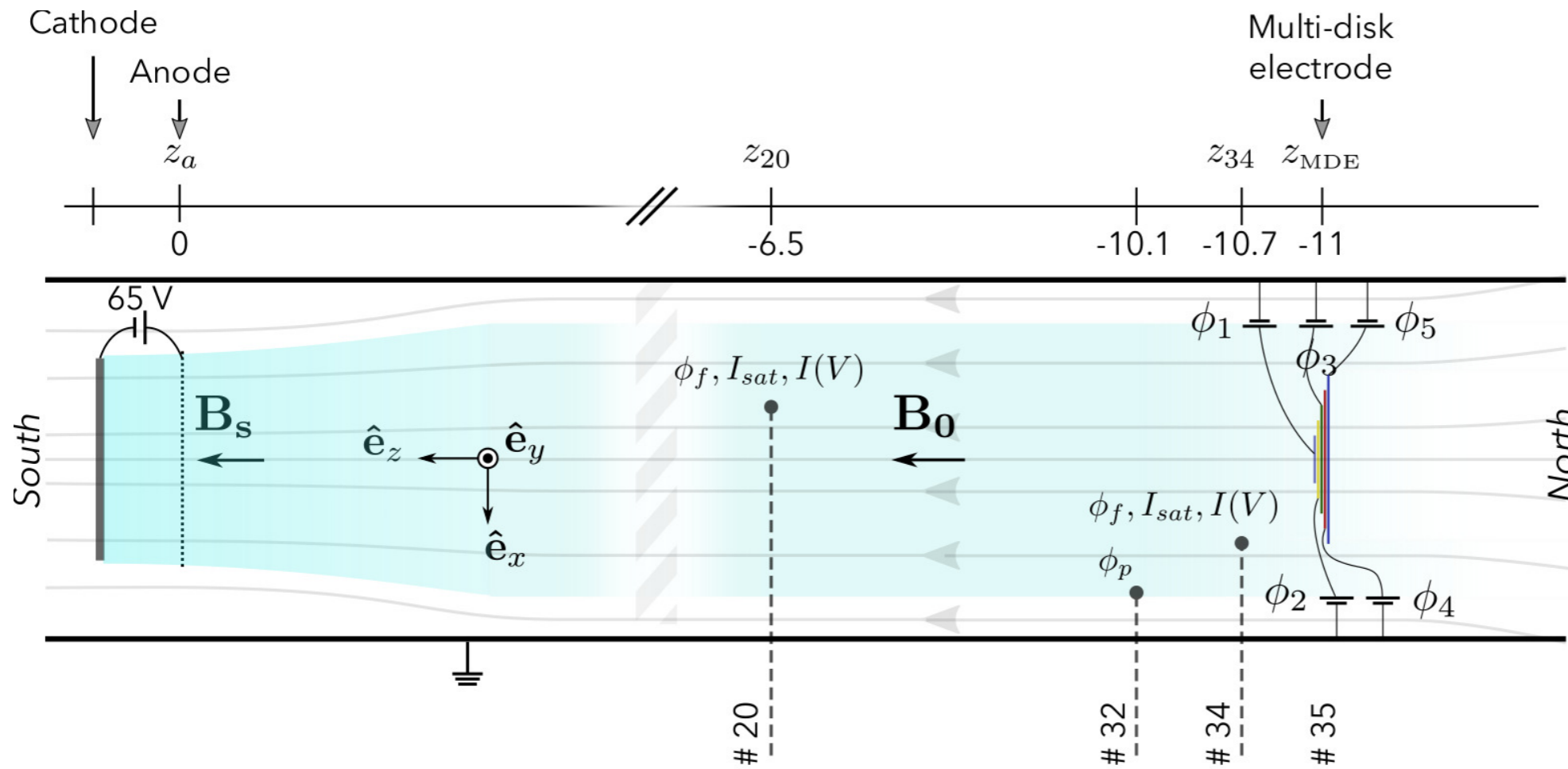


UCLA Physics and Astro Dpt - US NSF/DOE User Facility

- Magnetic field up to 2.5 kG uniform over 20 m.
- Plasma source: hot cathode. LaB6 disk \varnothing 38 cm.
- Pulsed operation, typical repetition rate of 1 Hz.
- $n_e \leq \text{few } 10^{19} \text{ m}^{-3}$, $T_e, T_i \leq 20 \text{ eV}$, partially to nearly fully ionized.
- Access ports for automated probe diagnostics every 30 cm.

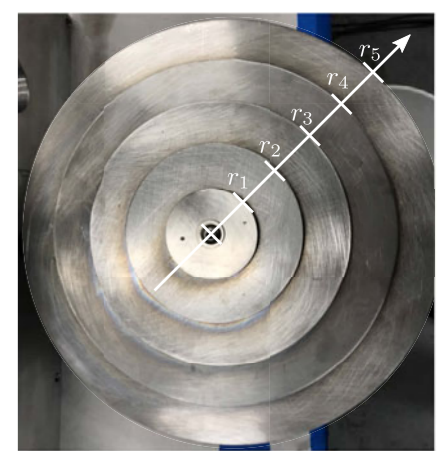


Objective: study the effect of end-electrode biasing - towards rotation control

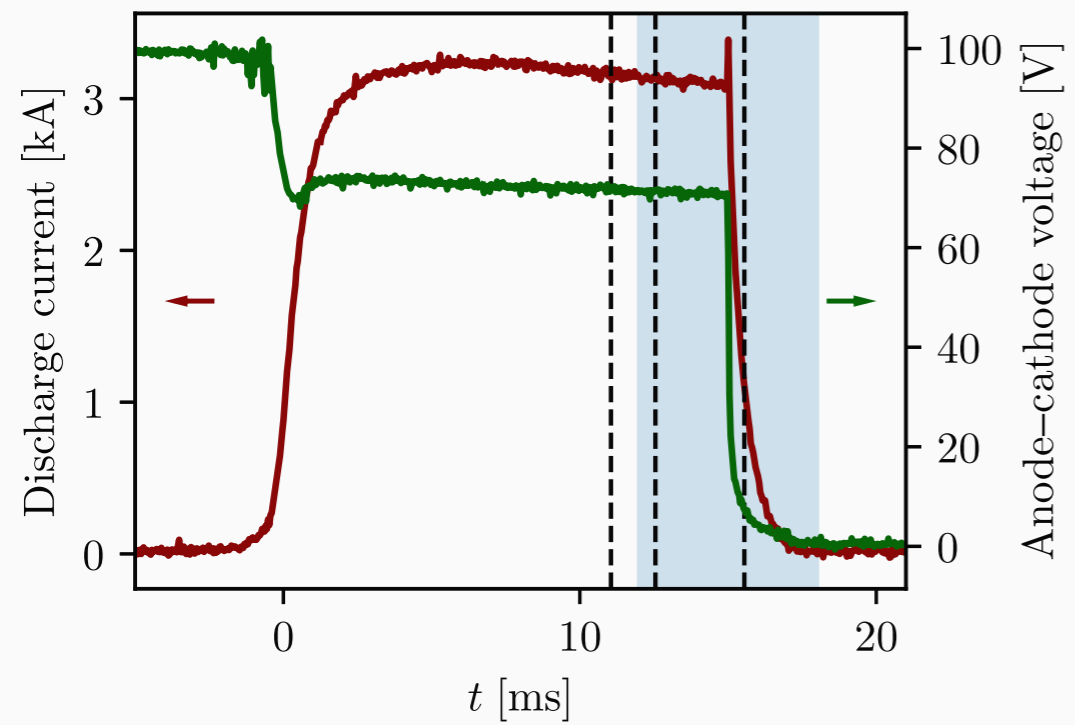


$r_i \sim 2.5i$ cm

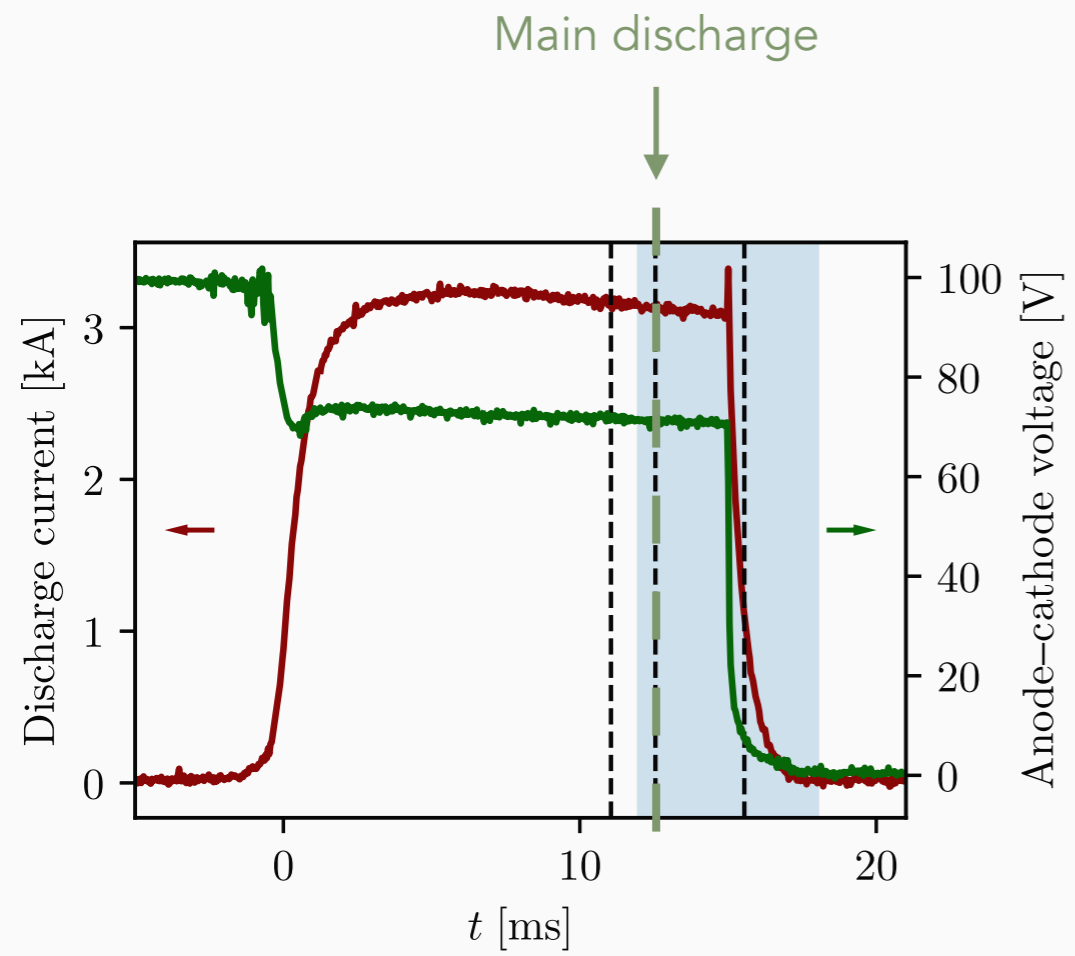
Parameter	Value
Gas	helium
Discharge current I_d [kA]	3.2
Main chamber magnetic field B_0 [T]	0.1
Magnetic field ratio ρ_b	2
Plasma density n_e [m^{-3}]	$5 \cdot 10^{18}$
Electron temperature T_e [eV]	5
Plasma radius [cm]	25
Plasma length [m]	11



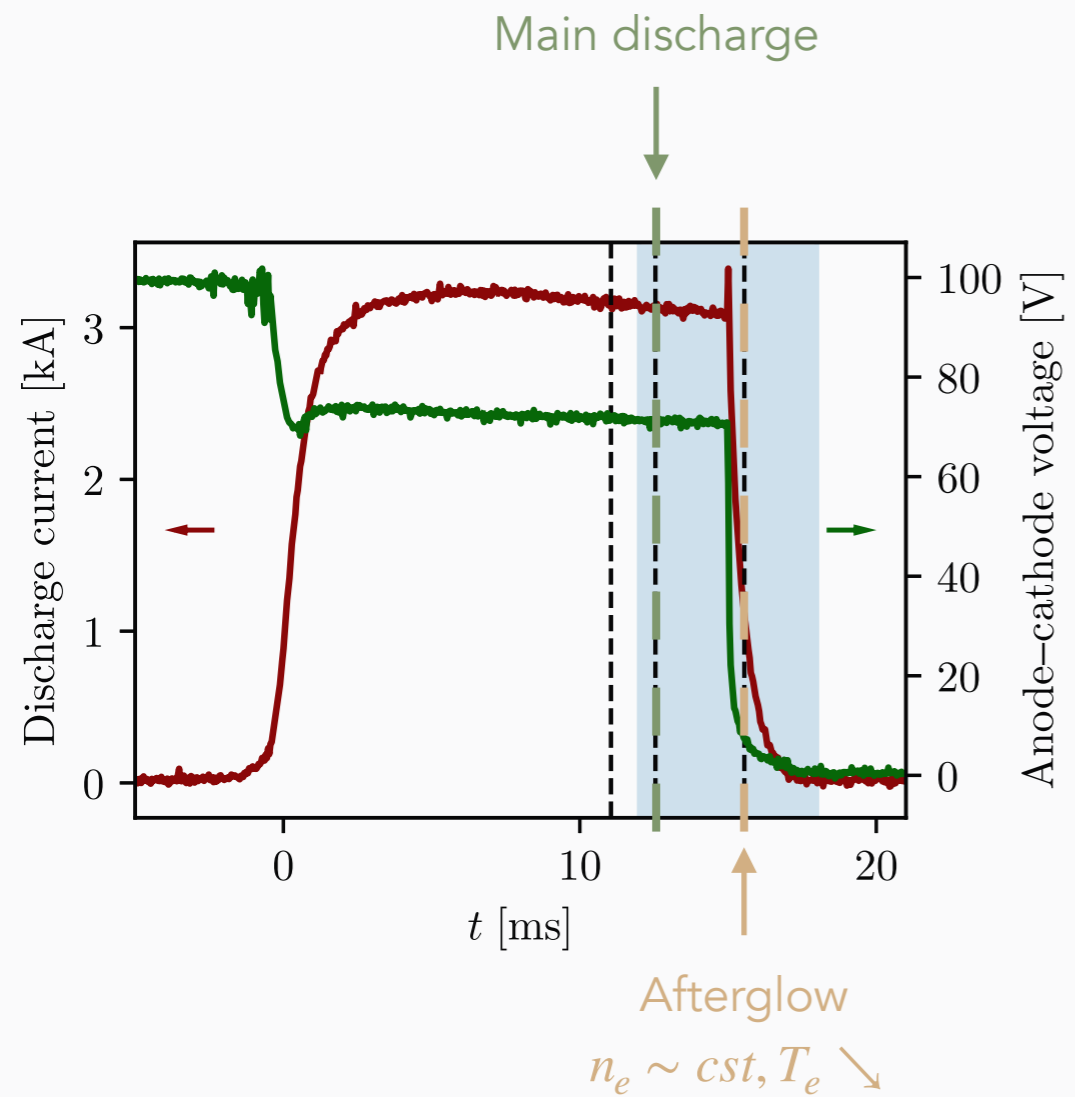
Data at different instants during a shot



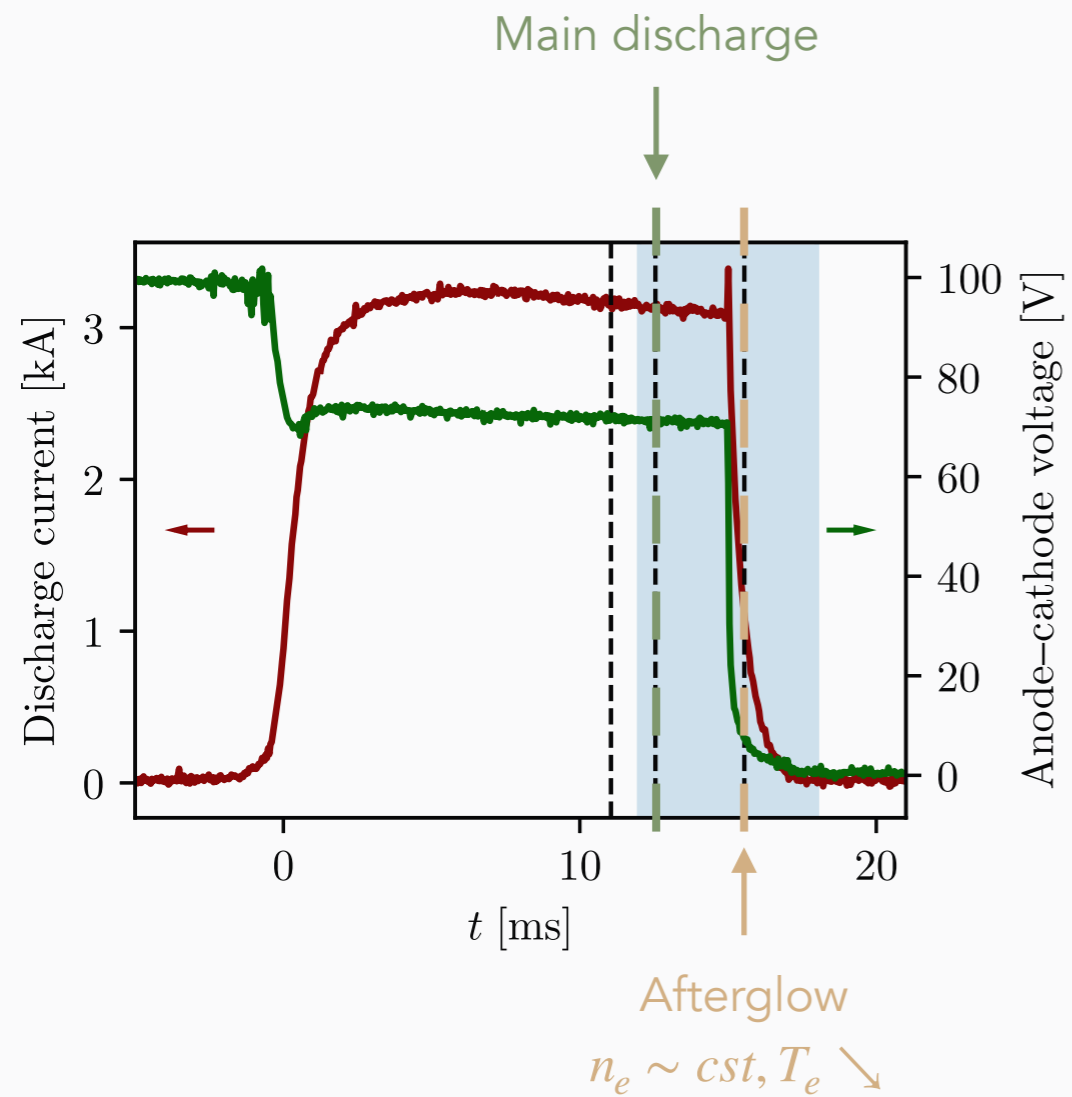
Data at different instants during a shot



Data at different instants during a shot

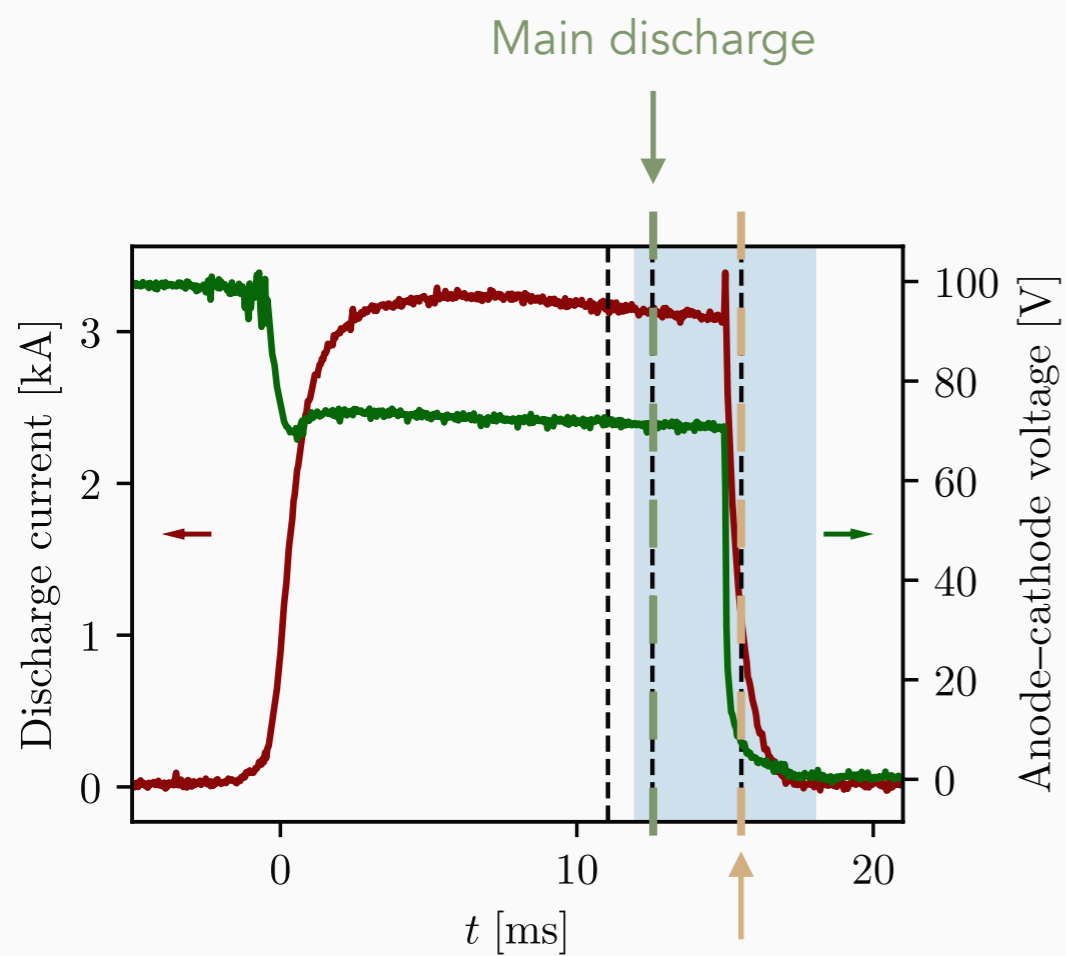


Data at different instants during a shot



Today: "main discharge" only

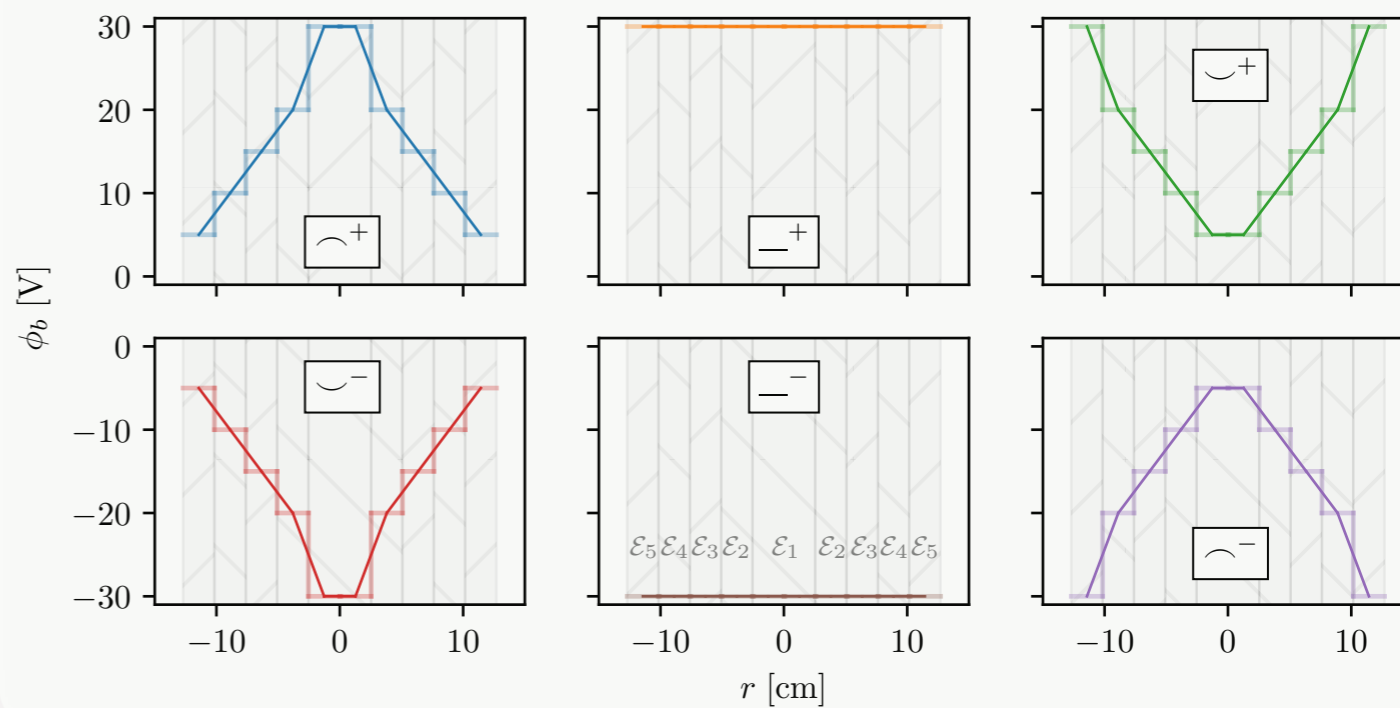
Data at different instants during a shot

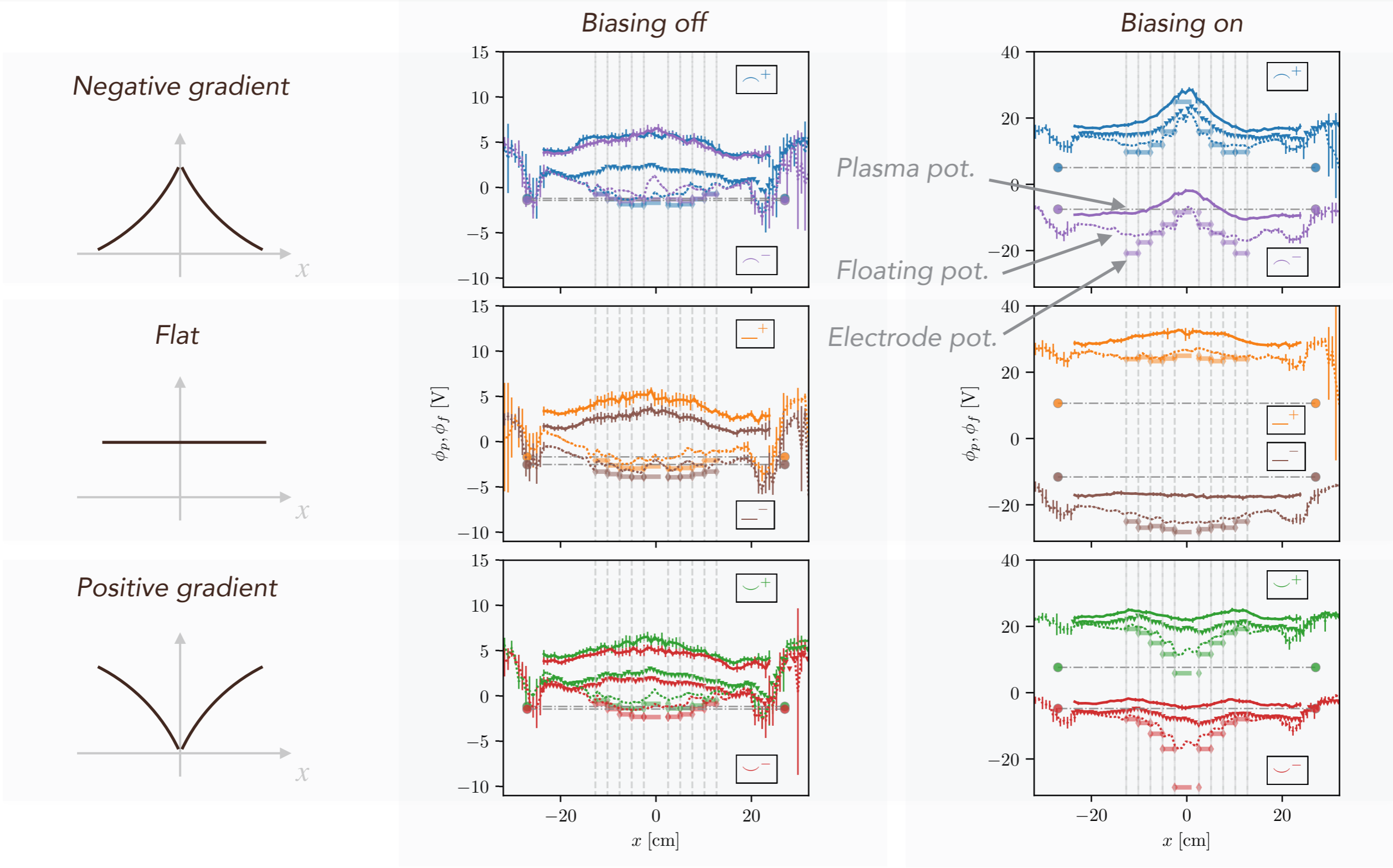


Afterglow
 $n_e \sim cst, T_e \searrow$

Today: "main discharge" only

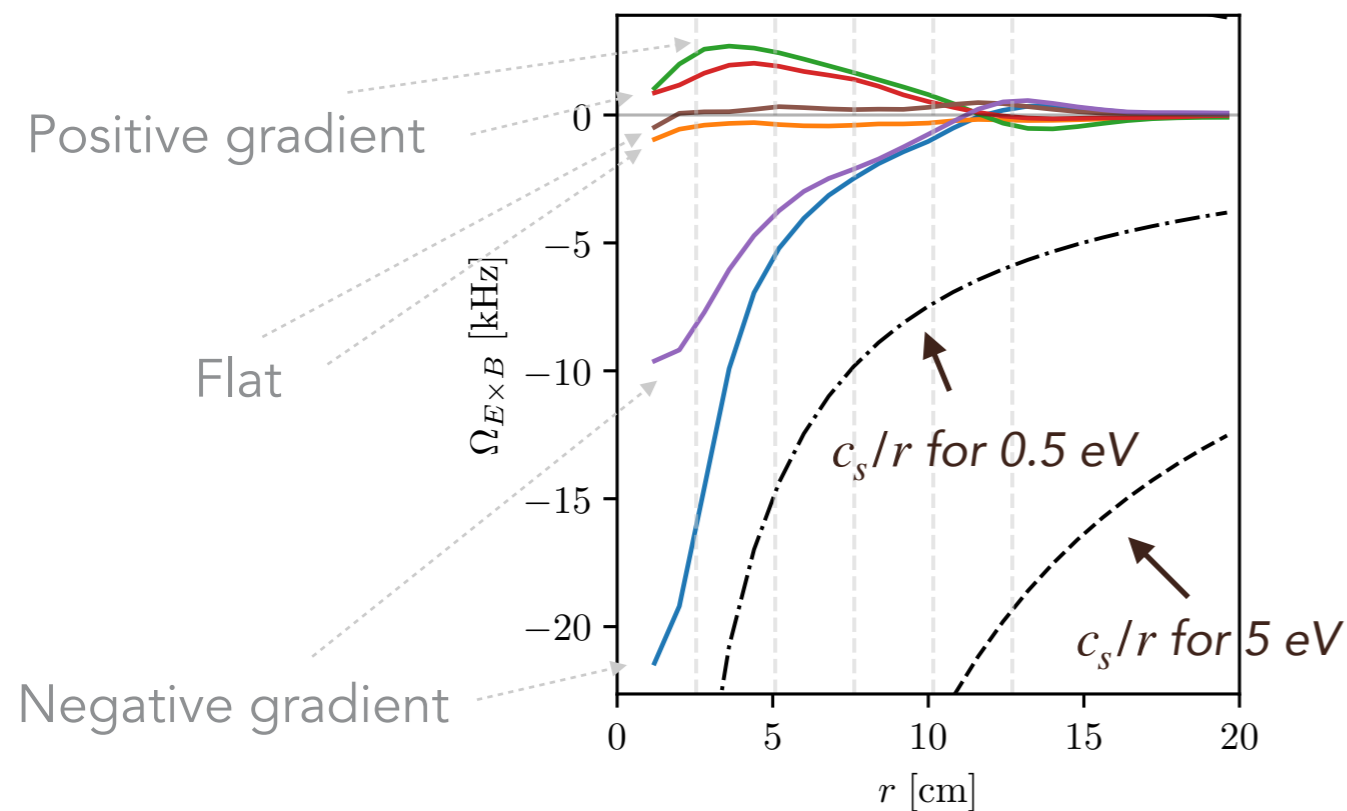
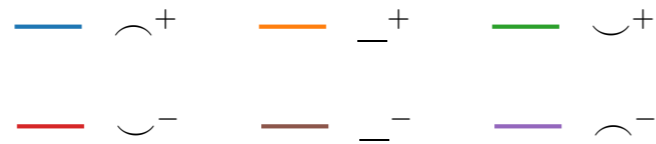
6 biasing scenarios, both polarities, both potential gradient signs





Biasing has a significant effect on the plasma potential in the machine:

- Amplitude comparable to applied bias (± 30 V)
- Recovered 6 meters away from electrodes



Relatively large (but subsonic) rotation

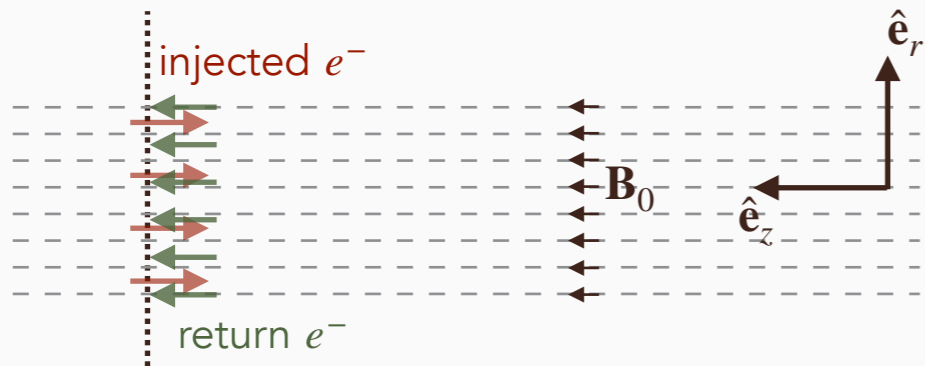
- near the axis for negative gradients imposed on the electrodes

Data for other biasing scenarios, and more generally for slower rotation, should be compared to diamagnetic contributions.

The anode potential is not fixed. It self adjusts to maintain overall current balance.

Without electrodes

Floating
mesh anode



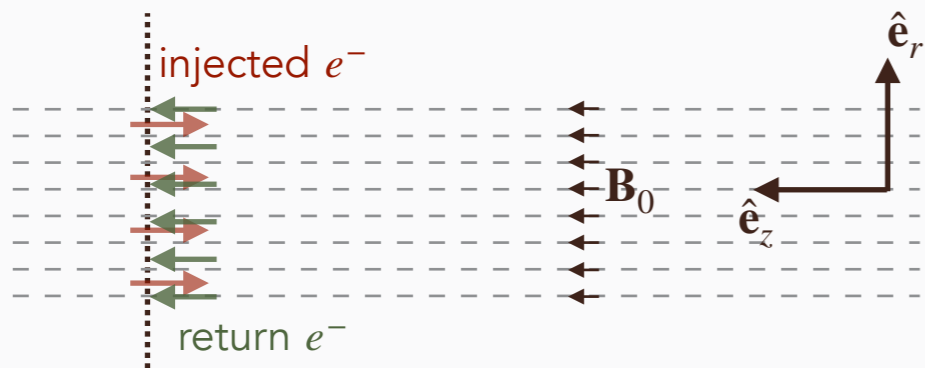
$$\int j_{inj} dS = \int j dS$$

The anode seats below the local plasma potential to draw an electron current

The anode potential is not fixed. It self adjusts to maintain overall current balance.

Without electrodes

Floating
mesh anode

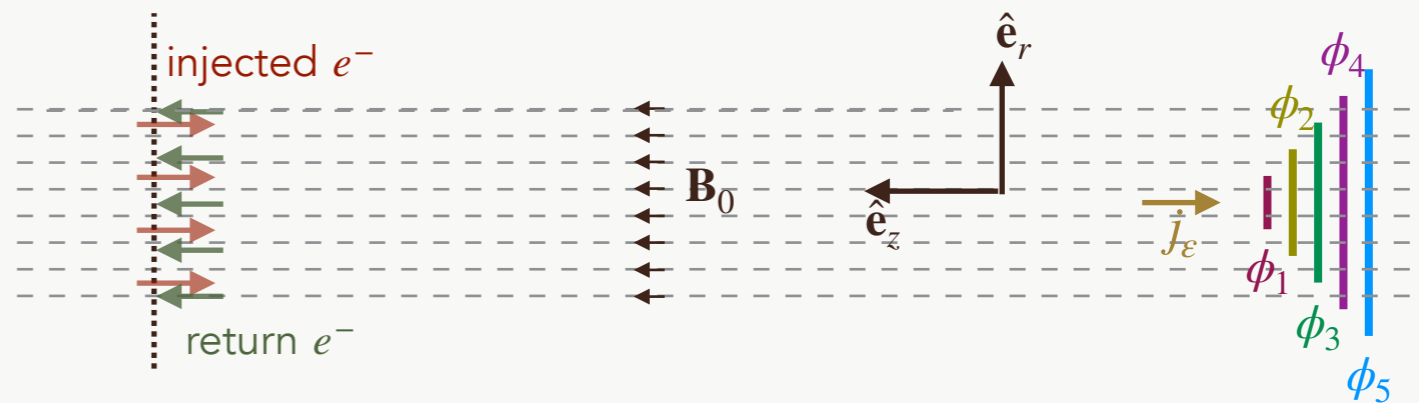


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With electrodes

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mesh anode



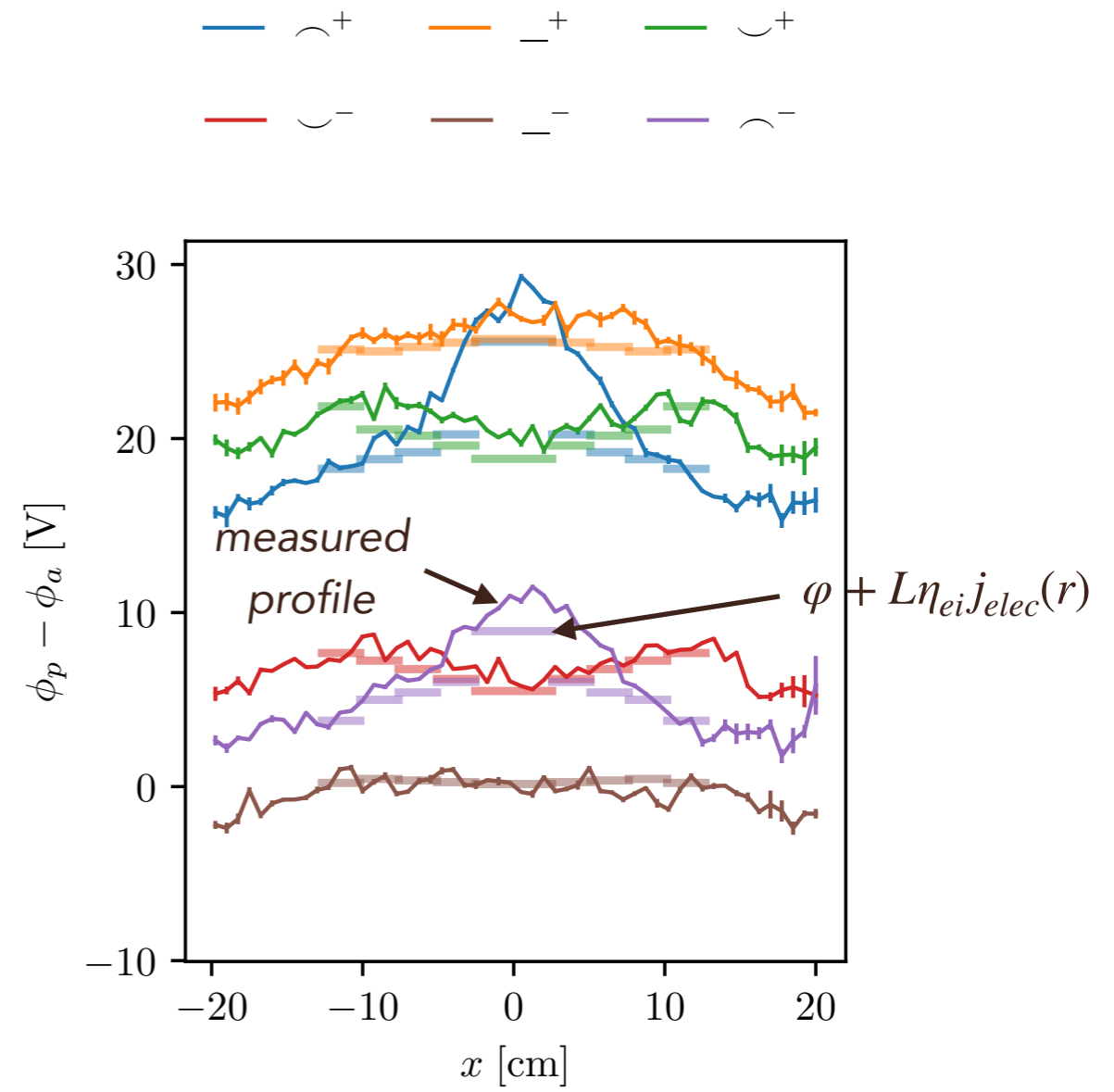
$$\int j_{\epsilon} dS_{\epsilon} \neq 0 \rightarrow \int j_{inj} dS \neq \int j dS$$

Biased electrodes draw net current which offsets the balance at the anode → affects the “anode sheath”
Net current also leads to axial voltage drop

Potential radial profile is well approximated by the axial voltage drop computed from the Spitzer resistivity and the current collected on each electrode.

$$\phi_p(r) = \phi_a + \varphi + L\eta_{ei}j_{elec}(r)$$

anode pot. (measured) axial resistive drop (inferred from plasma n_e, T_e and current on electrodes)
 plasma pot. (measured) anode sheath (free parameter)

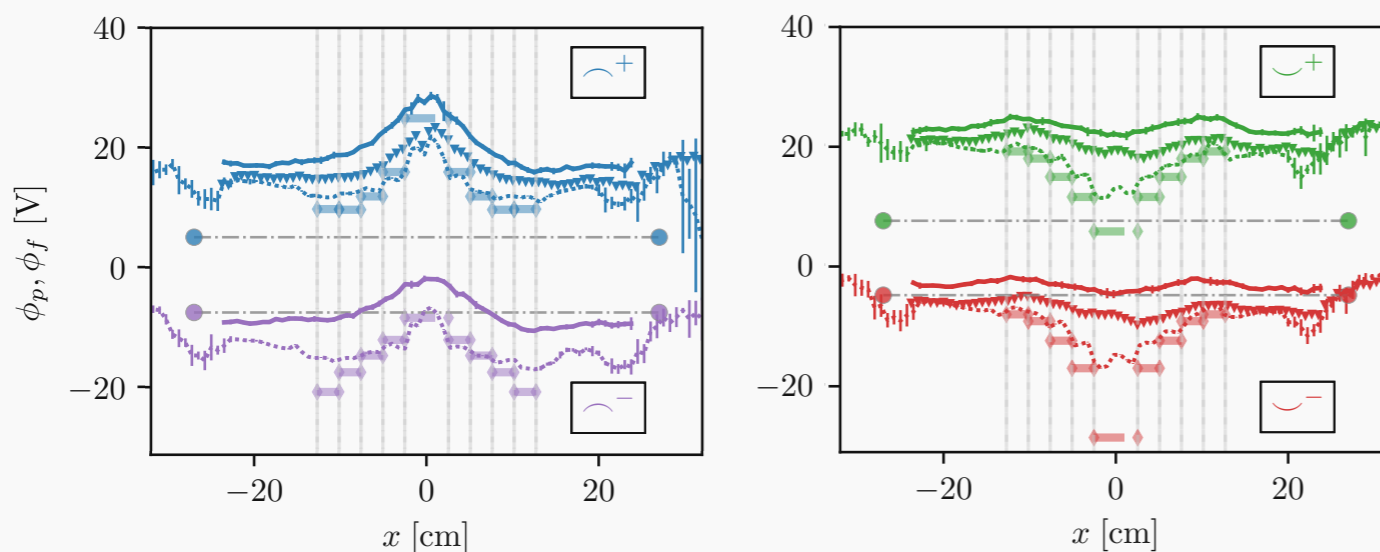


Radial potential profile is driven by the axial current drawn on biased electrodes.

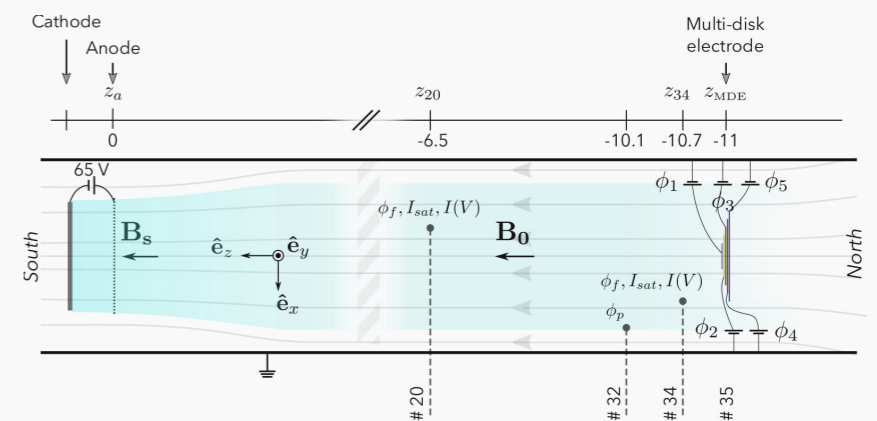
Gueroult et al. (2024), arXiv:2401.06480

- ▶ Controlling plasma rotation has high upside potential, both for basic physics and for applications: wave physics, fusion, separation, ...
- ▶ One long proposed solution is to use biased electrodes... but experimental results so far have been mixed. Likely causes are the sheath and the voltage drops along field line at finite $\sigma_{\perp}/\sigma_{\parallel}$.
- ▶ New results obtained in LAPD show that the radial electric field can indeed be affected/controlled via axial currents and finite $\sigma_{\perp}/\sigma_{\parallel}$, but many questions remain including:

Asymmetry in positive and negative gradients

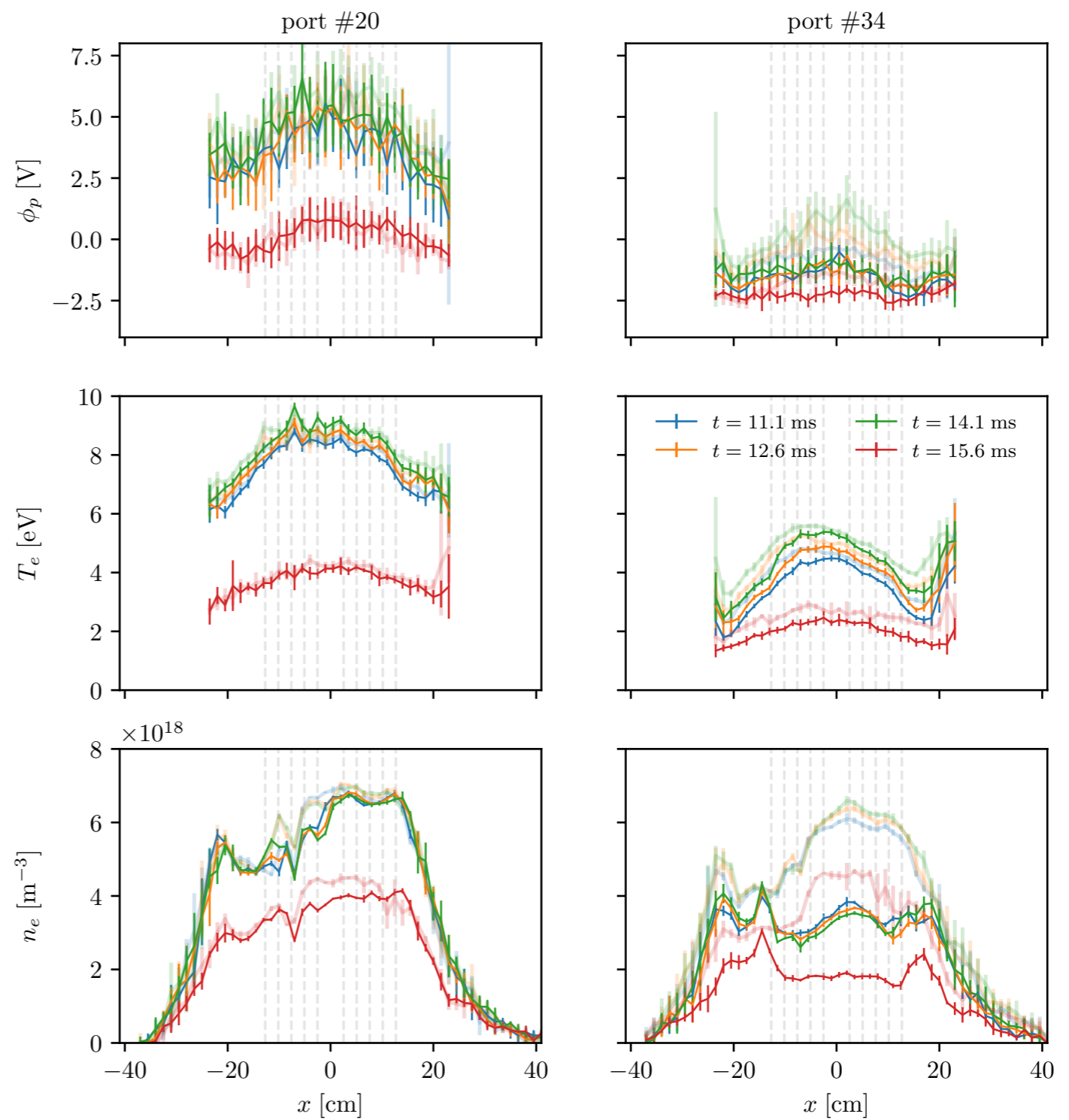


Symmetric end conditions

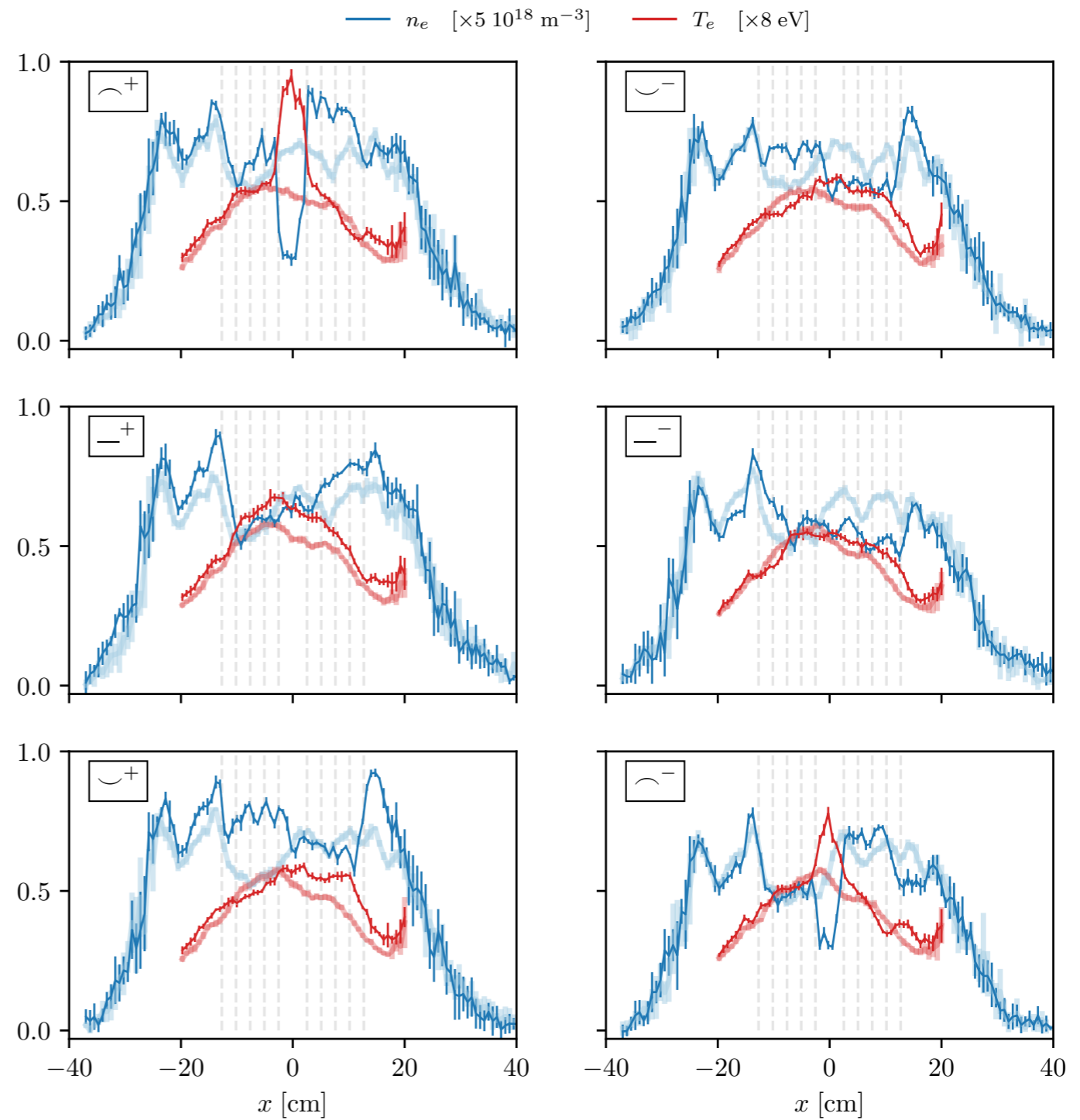


EXTRA SLIDES

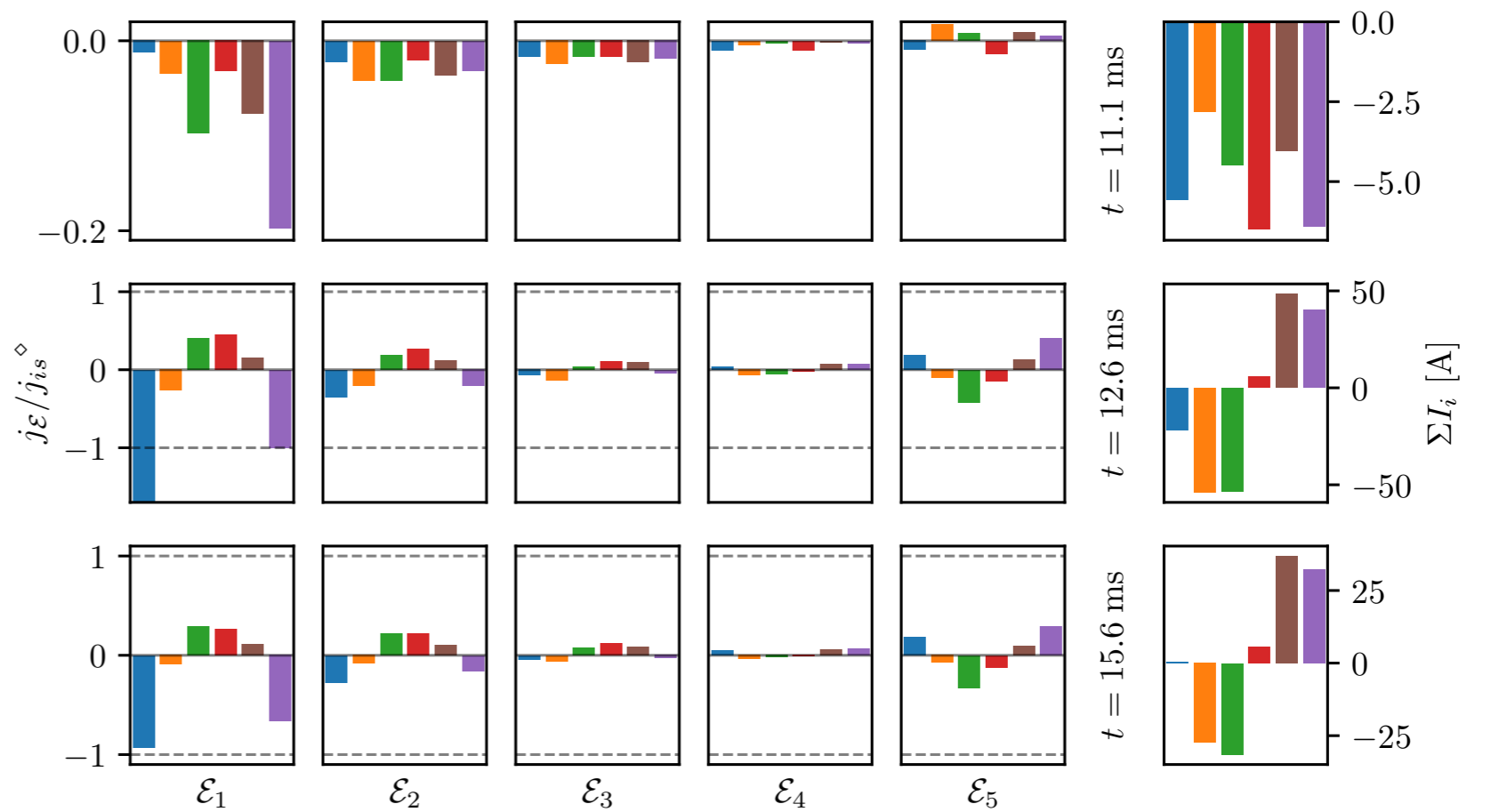
Radial profiles at different instants with and without the electrodes in the machine



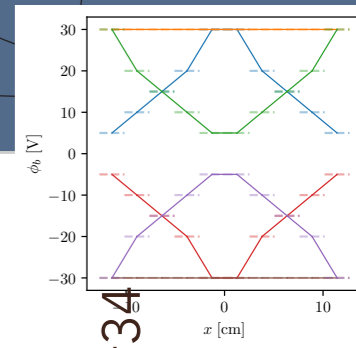
Radial profiles of density and temperature before and during active biasing



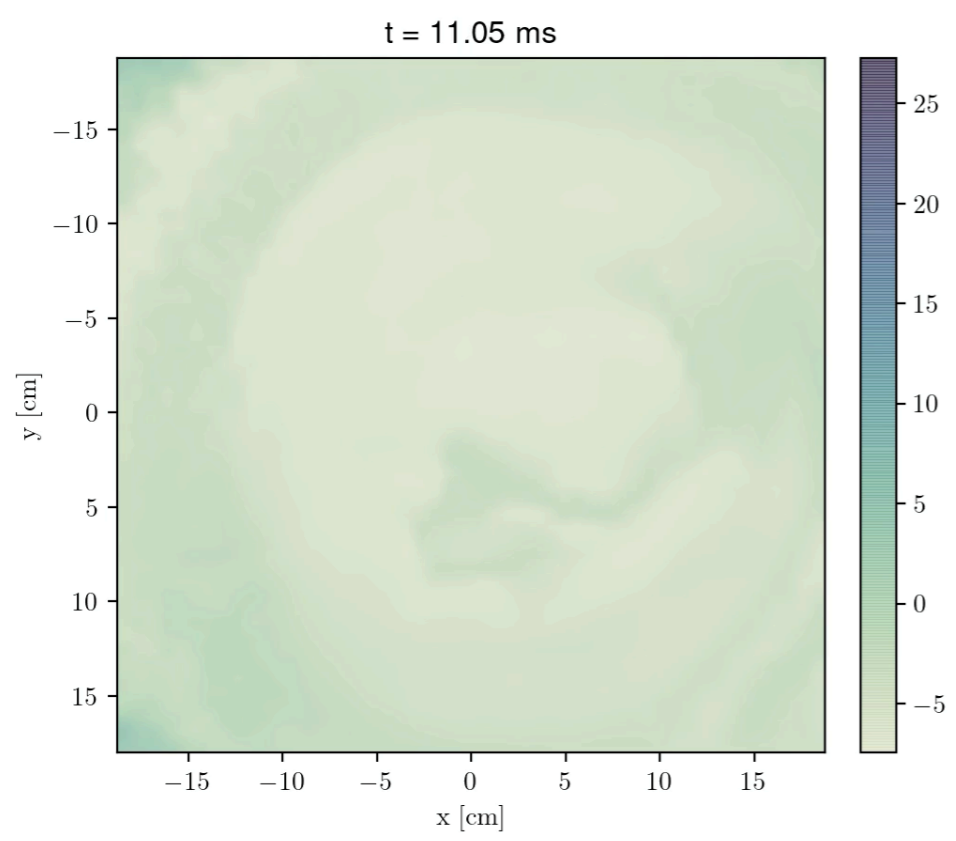
Current on electrodes



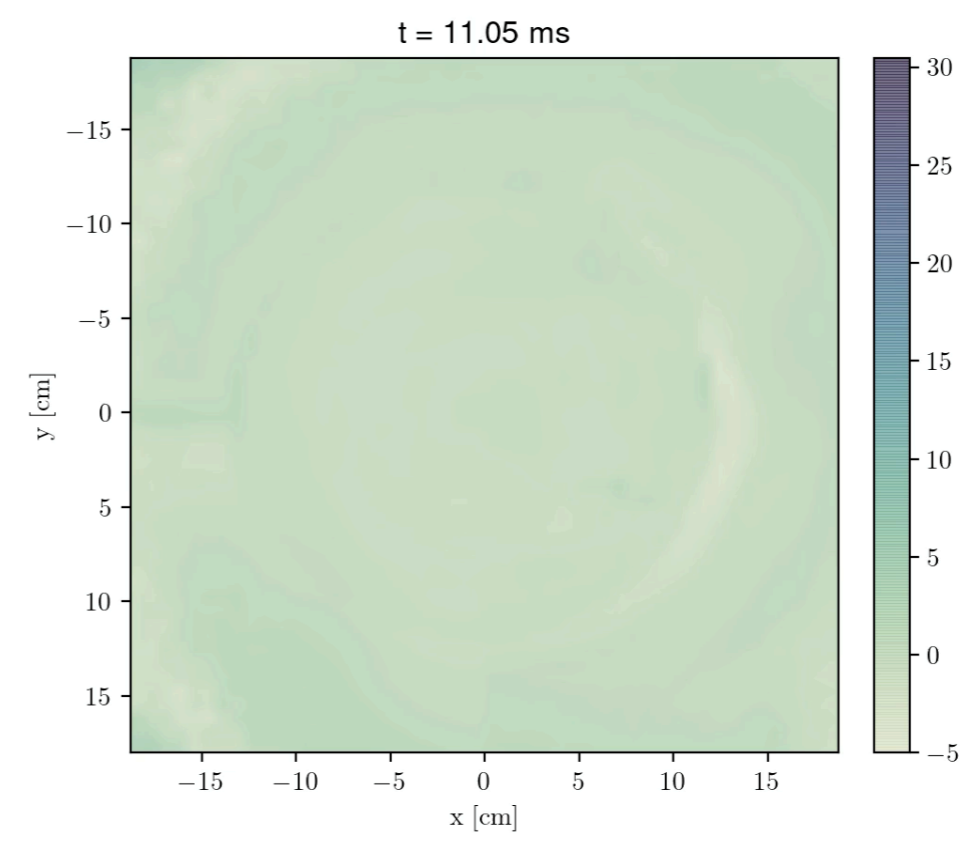
2D MAPS: +30 V INVERTED = GREEN



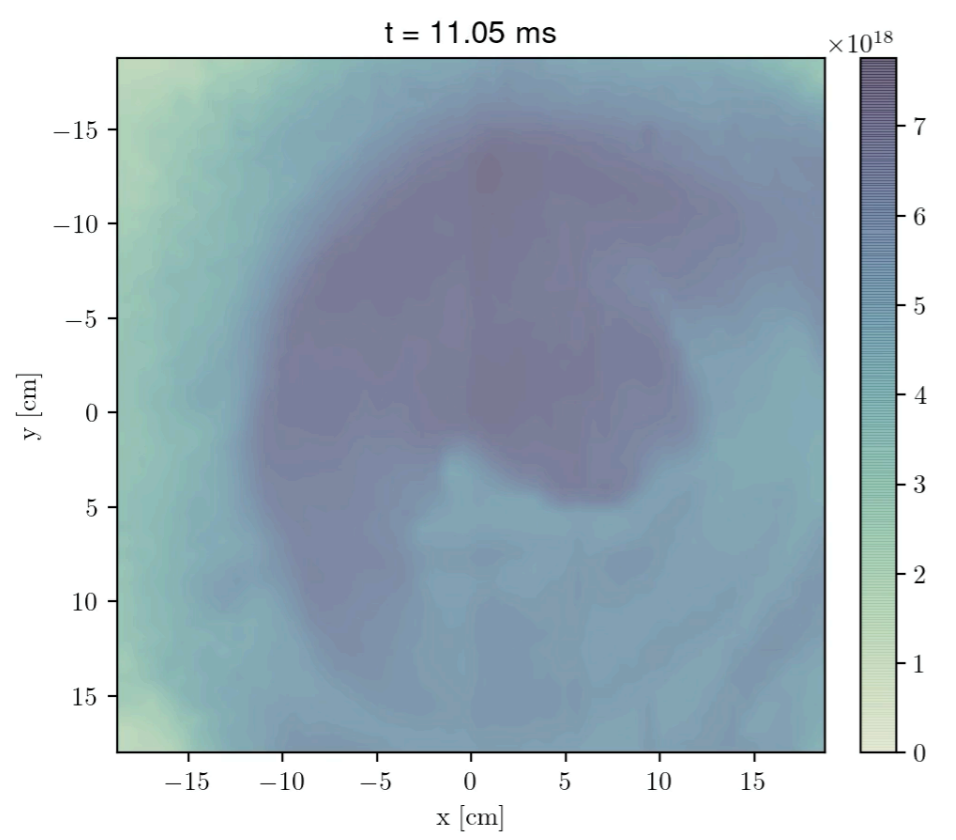
Floating potential port #20



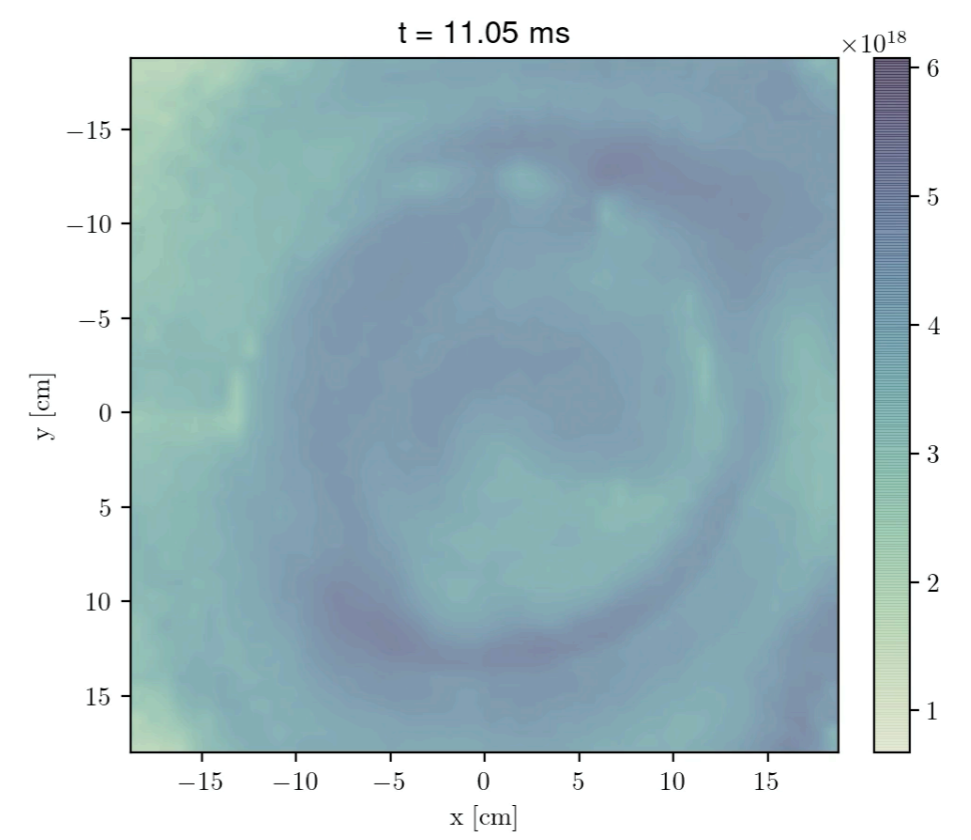
Floating potential port #34



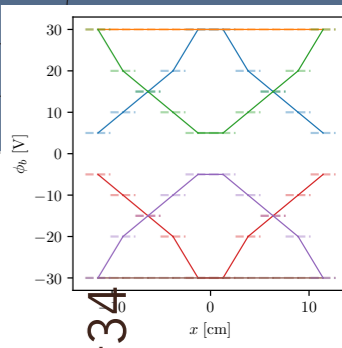
Density port #20



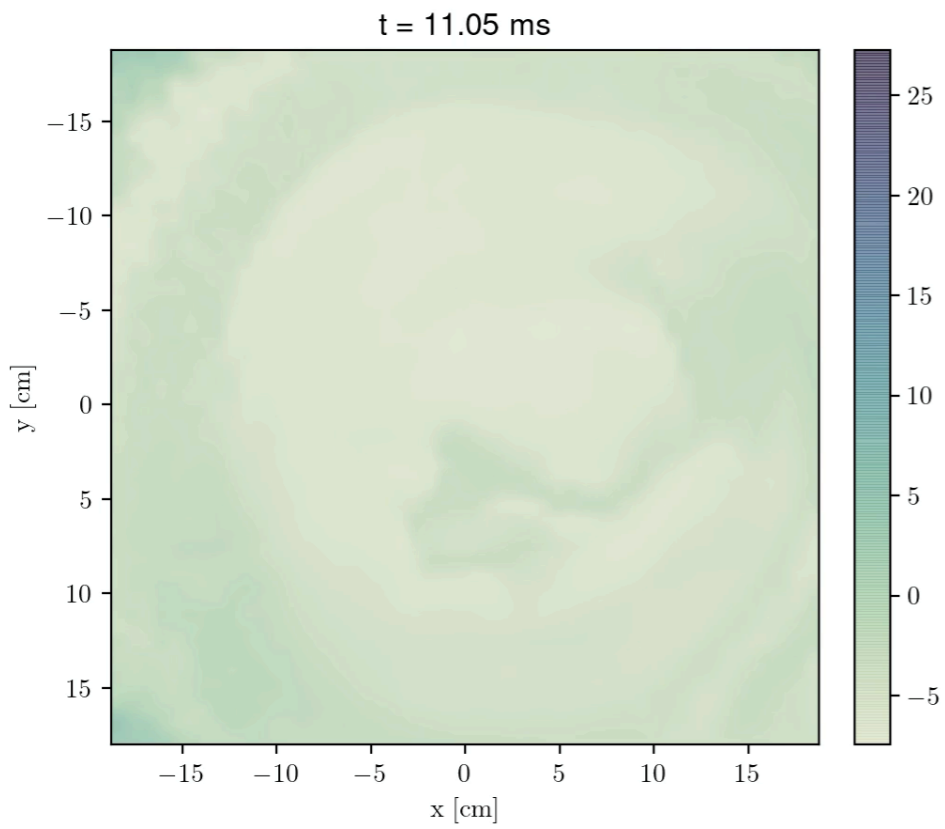
Density port #34



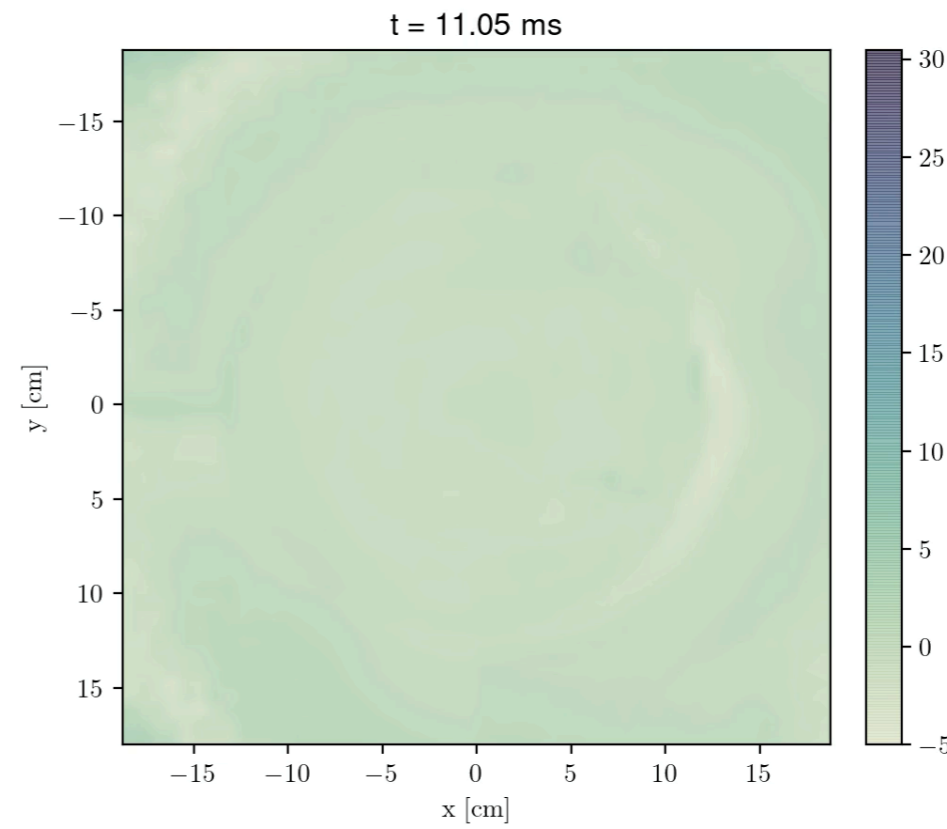
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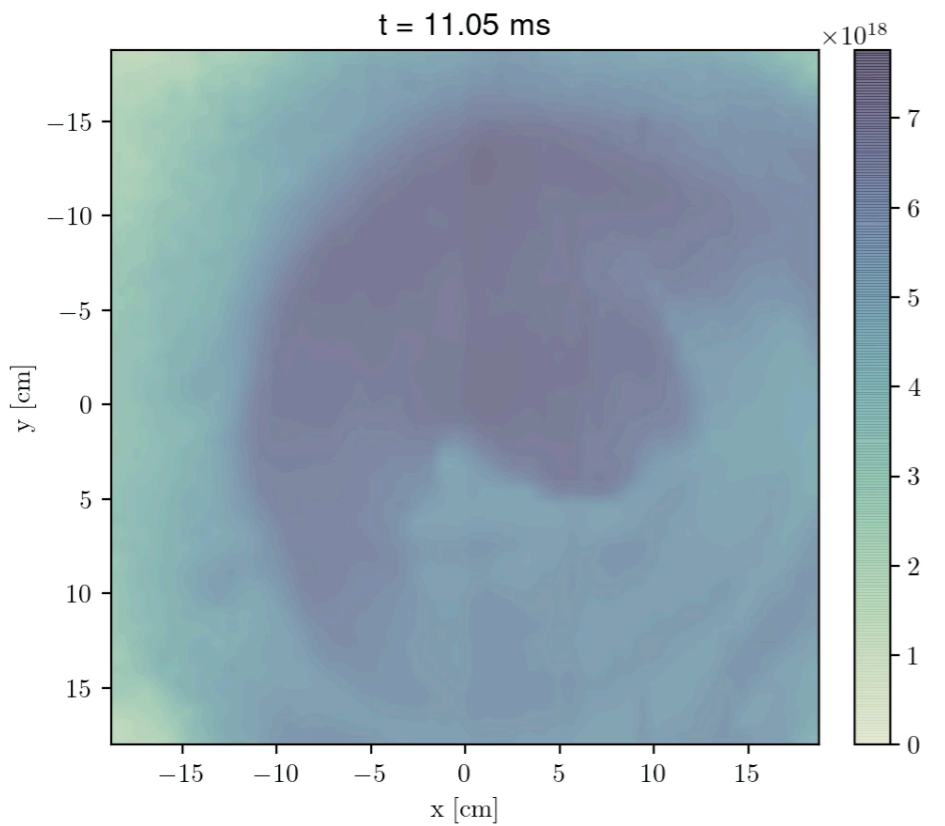
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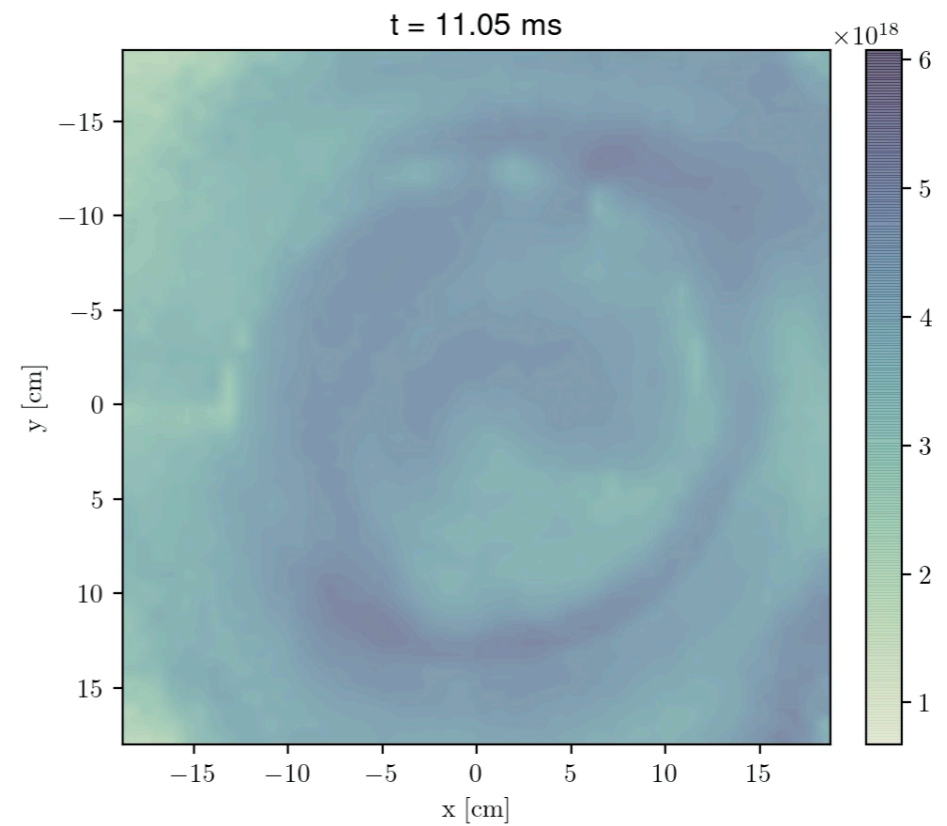
Floating potential port #34



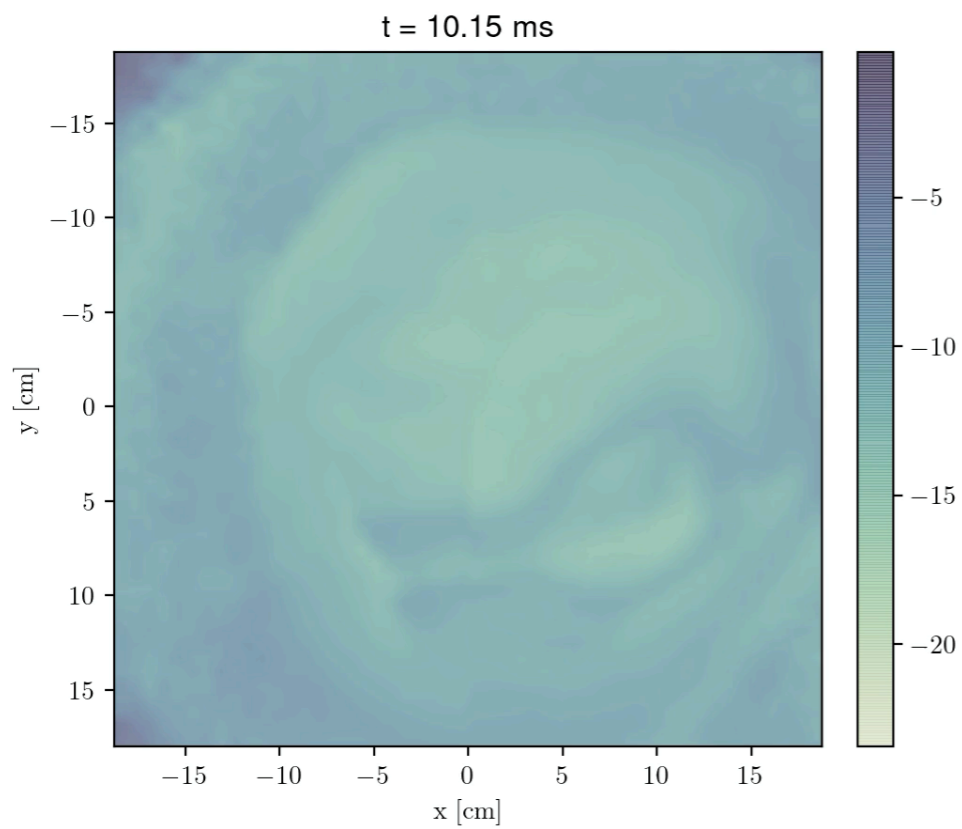
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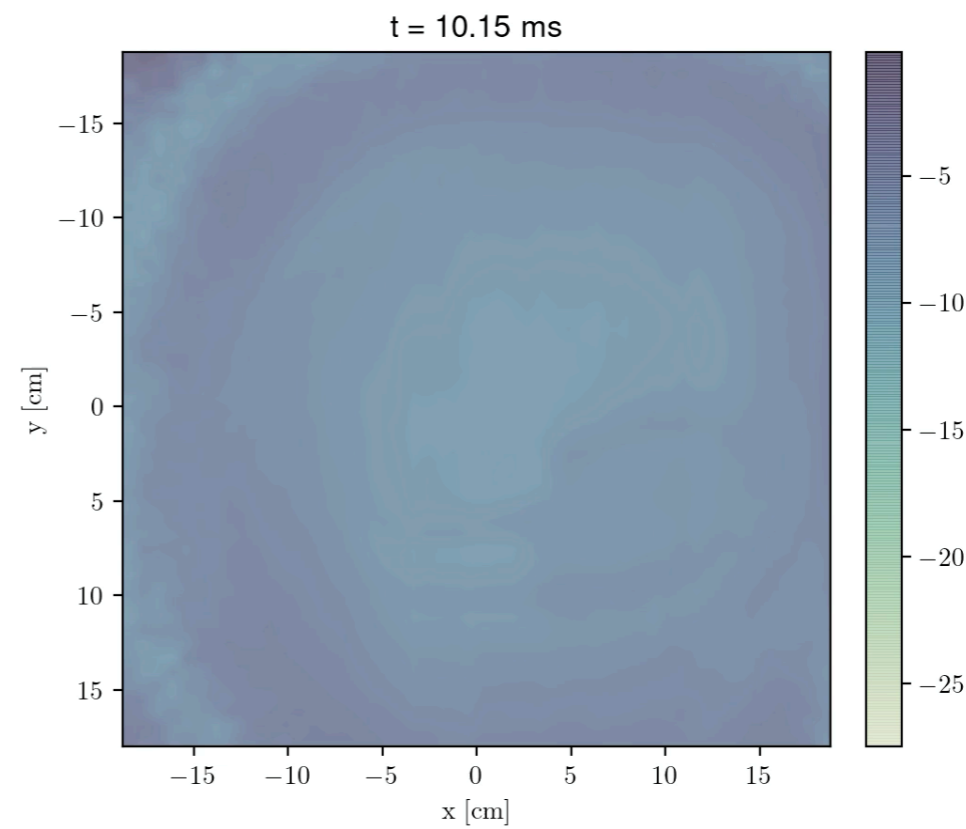
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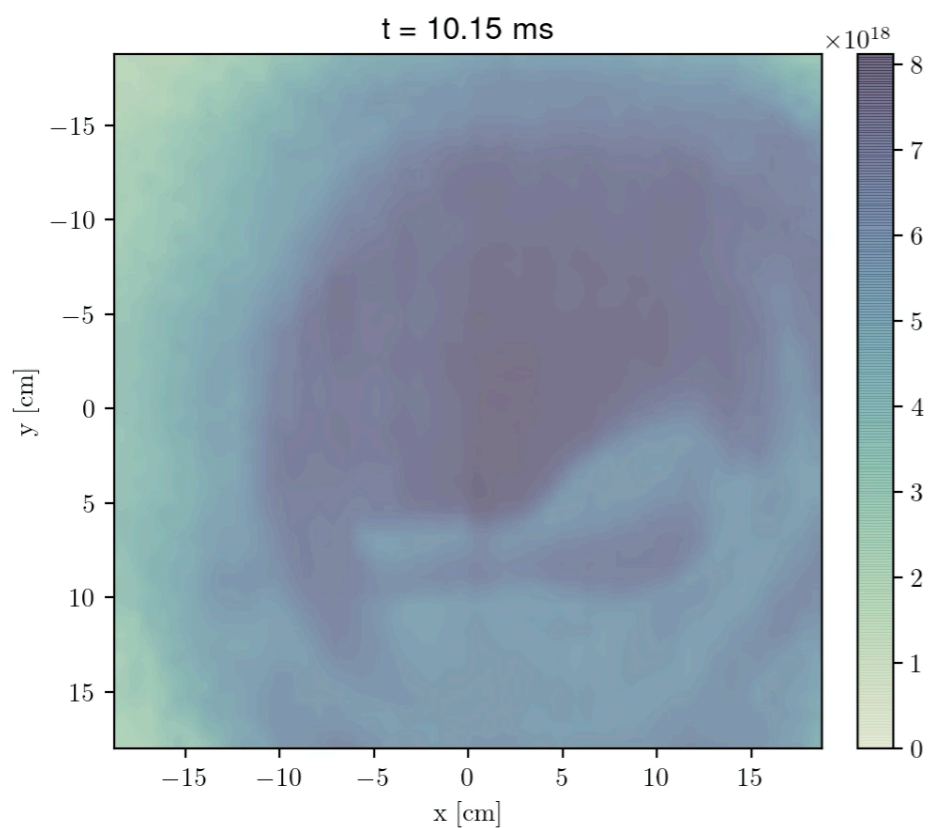
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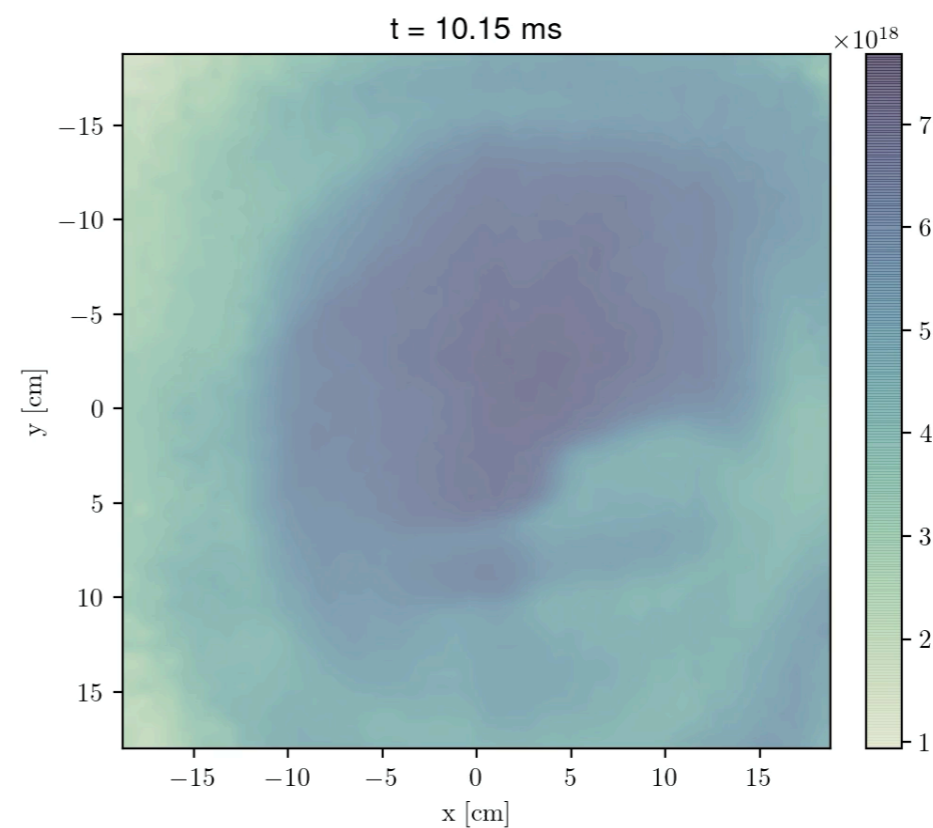
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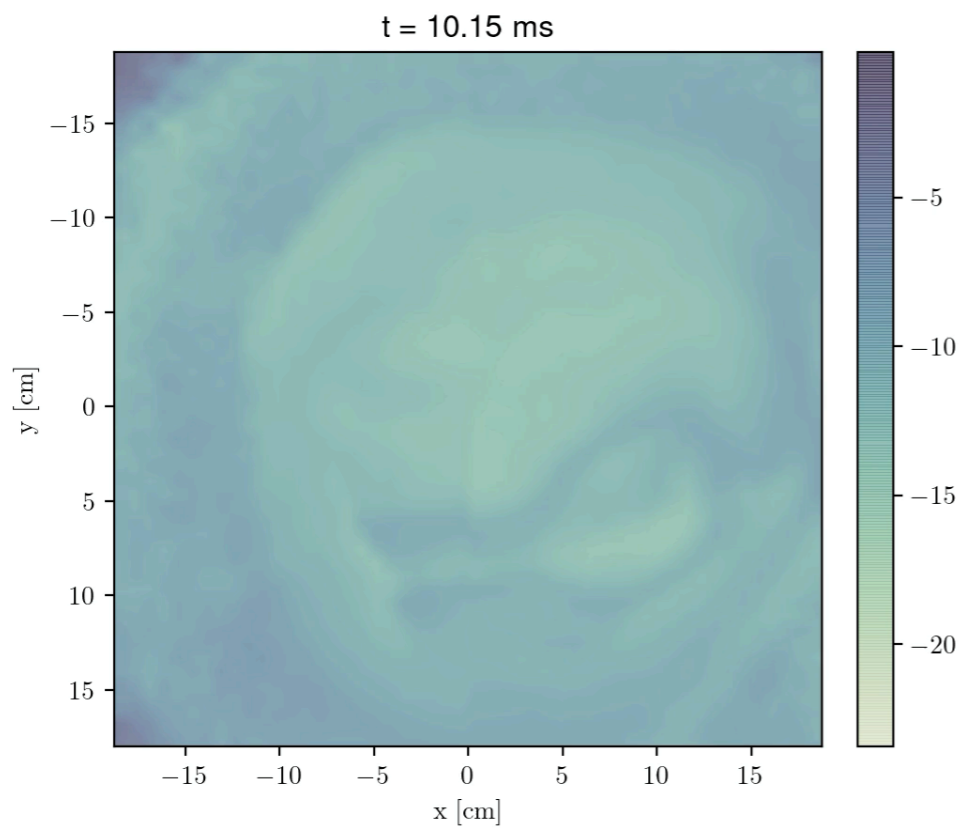
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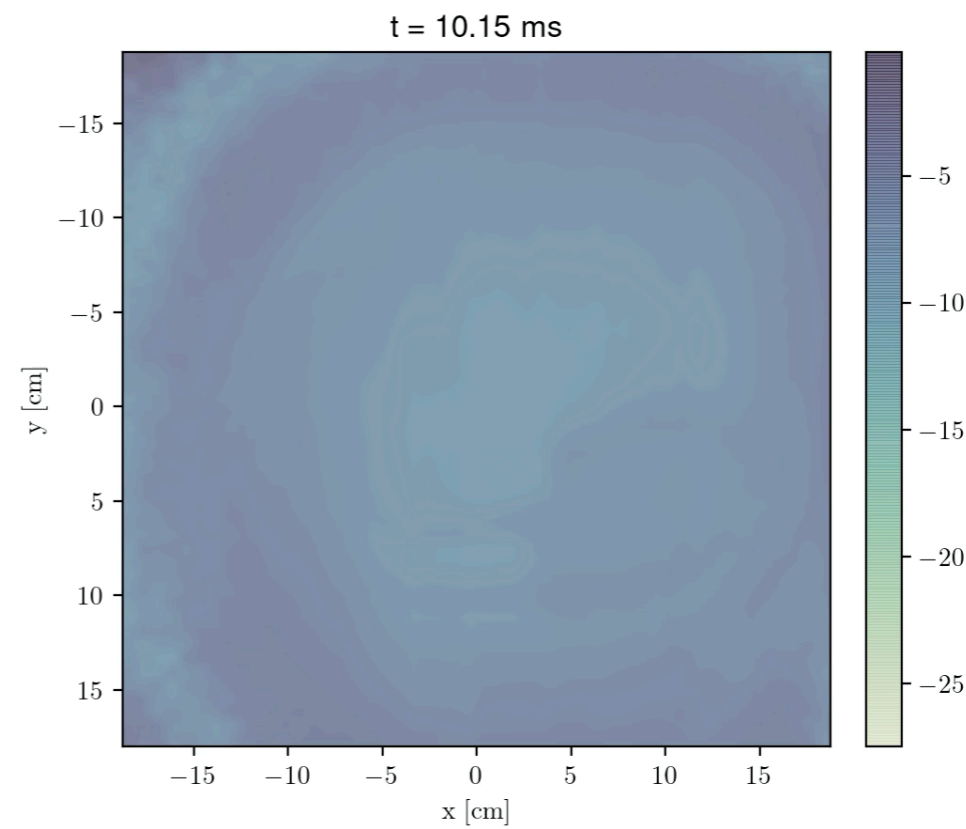
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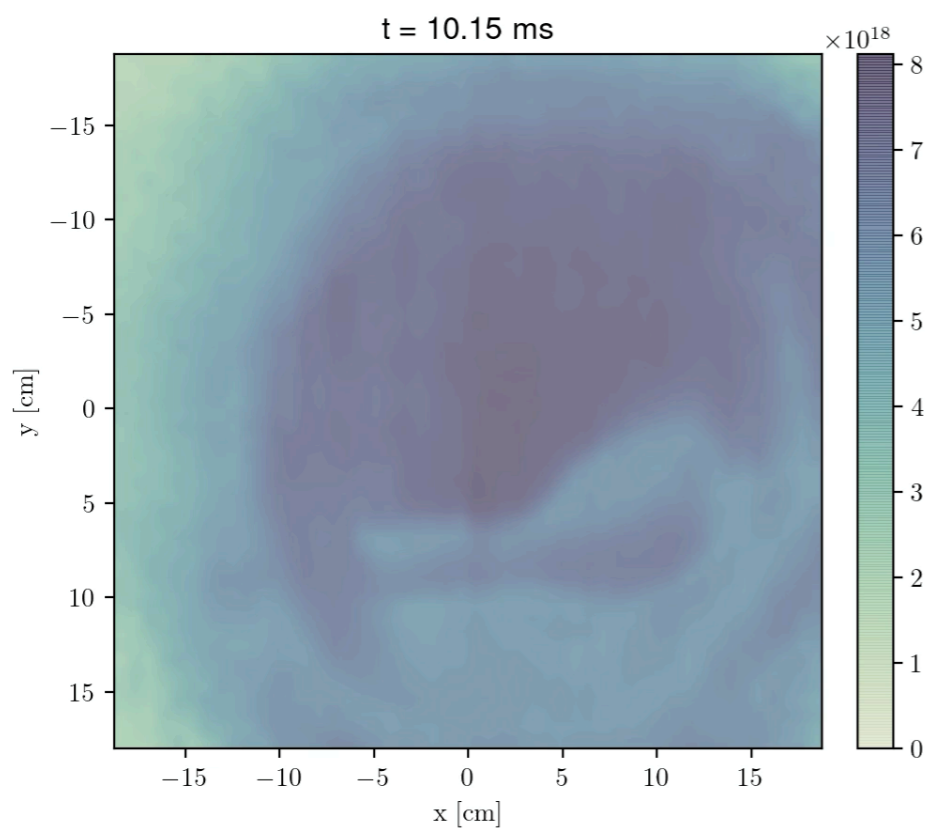
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